





Rheology of the Mantle and Plates (part 1): Deformation mechanisms and flow rules of mantle minerals





Topics covered in this class

- Rheology of Earth (viscous limit)
- Fluid Dynamics for geological phenomena
- Composition of the Earth
- Thermodynamics and high-pressure mineral physics
- Seismological structure of the mantle
- Geochemical structure of the mantle
- Dynamic processes of the Earth (plumes, slabs, thermochemical piles)
- Heat and mass transport in the deep Earth (convection, thermal history)
- Energetics of the core (magnetic field generation)



Where to turn to for more help..

- Mantle Convection in Earth + Planets
 Schubert, Turcotte, and Olson (2001)
- Numerical Geodynamic Modelling Gerya (2009)
- Hirth and Kohlstedt (2003)
- Regenaur-Lieb and Yuen (2003)
- Treatise on Geophysics V. 7, Ch. 2

- Papers by those lucky people in the Rheology fan club (partial list only):
 - numerical modelers: Podladchikov, Solomatov, Burov, Gerya, Bercovici, Tackley, Yuen
 - experimentalists: Karato, Kolhstedt, Hirth, Jackson



What is rheology?

- Rheology is the physical property that characterizes deformation behavior of a material (solid, fluid, etc)
- Rheology of Earth materials includes
 elasticity, viscosity, plasticity, etc
- For the deep Earth: mantle is fluid on geological timescales so we focus on its viscosity
- For tectonic plates: still viscous on geological timescales, but the effective viscosity is a subject of debate

$$\sigma=E\varepsilon~$$
 solid mechanics $\sigma=2\eta\dot{\varepsilon}~$ fluid mechanics

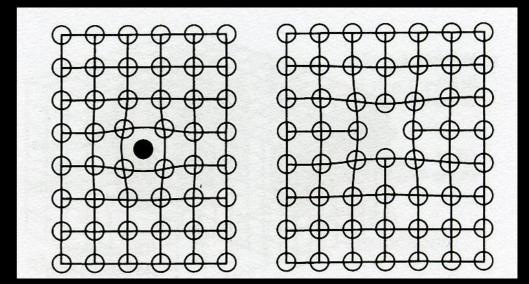


What is viscosity?

- constitutive relation between stress and strain-rate (deformation rate)
- in the continuum description, it is the analog of the elastic moduli which relate stress and strain
- measure of a fluid's ability to flow
- diffusivity of momentum

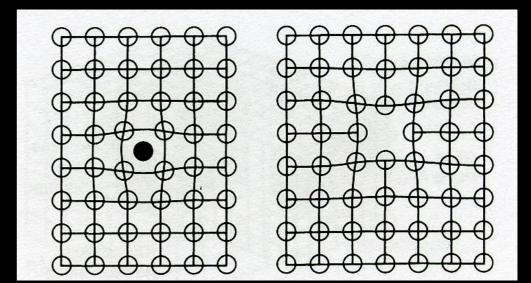


- movement of crystal defects
 - point defects extra atoms or vacancies



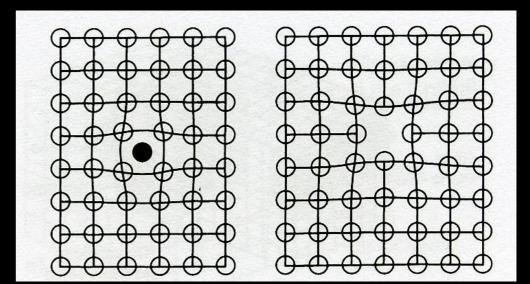


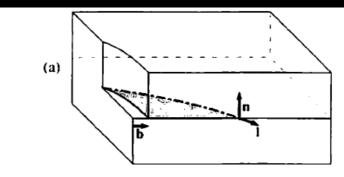
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 - line defects dislocations which represent a rearrangement of atomic bonds

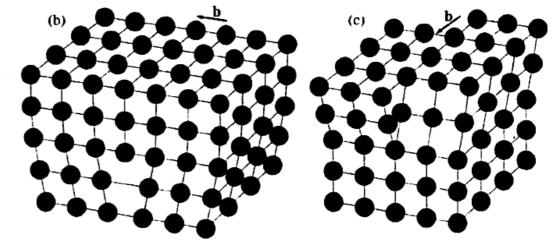




- movement of crystal defects
 - point defects extra atoms or vacancies
 - line defects dislocations which represent a rearrangement of atomic bonds
 - two types of dislocations: "edge" and "screw"
 - each described by parallel or normal Burgers vector (b*)

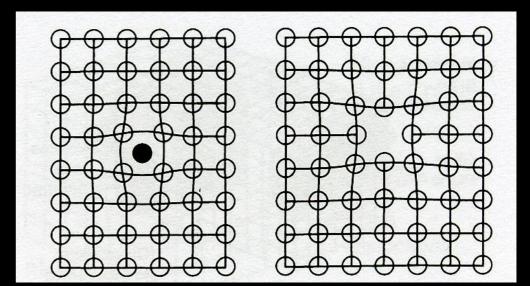


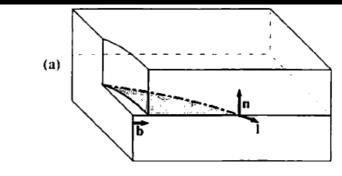


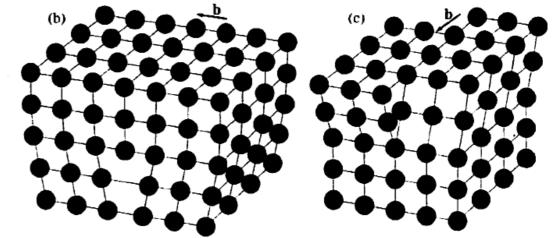




- movement of crystal defects
 - point defects extra atoms or vacancies
 - line defects dislocations which represent a rearrangement of atomic bonds
 - two types of dislocations: "edge" and "screw"
 - each described by parallel or normal Burgers vector (b*)
 - creep will occur through whichever mechanism requires least amount of energy





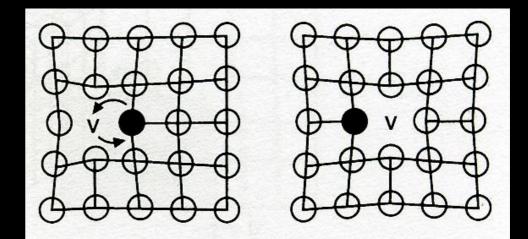




Diffusion creep

point defects move by diffusion

 through the crystal matrix (Nabarro-Herring)

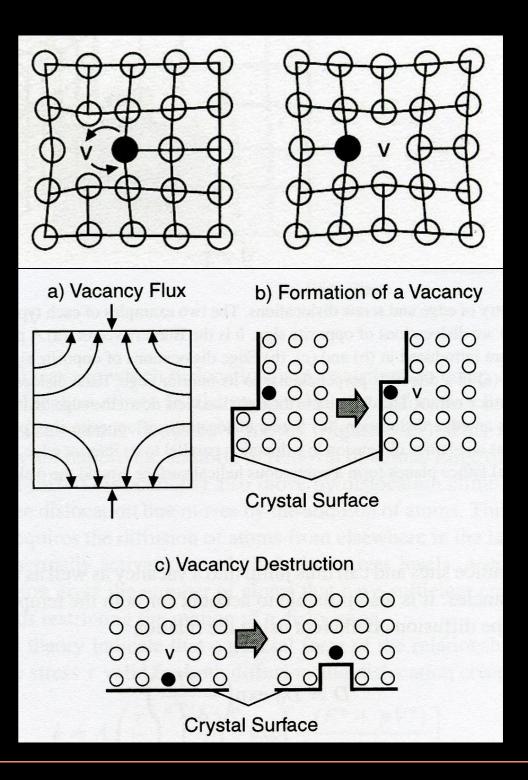


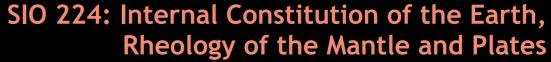


Diffusion creep

point defects move by *diffusion*

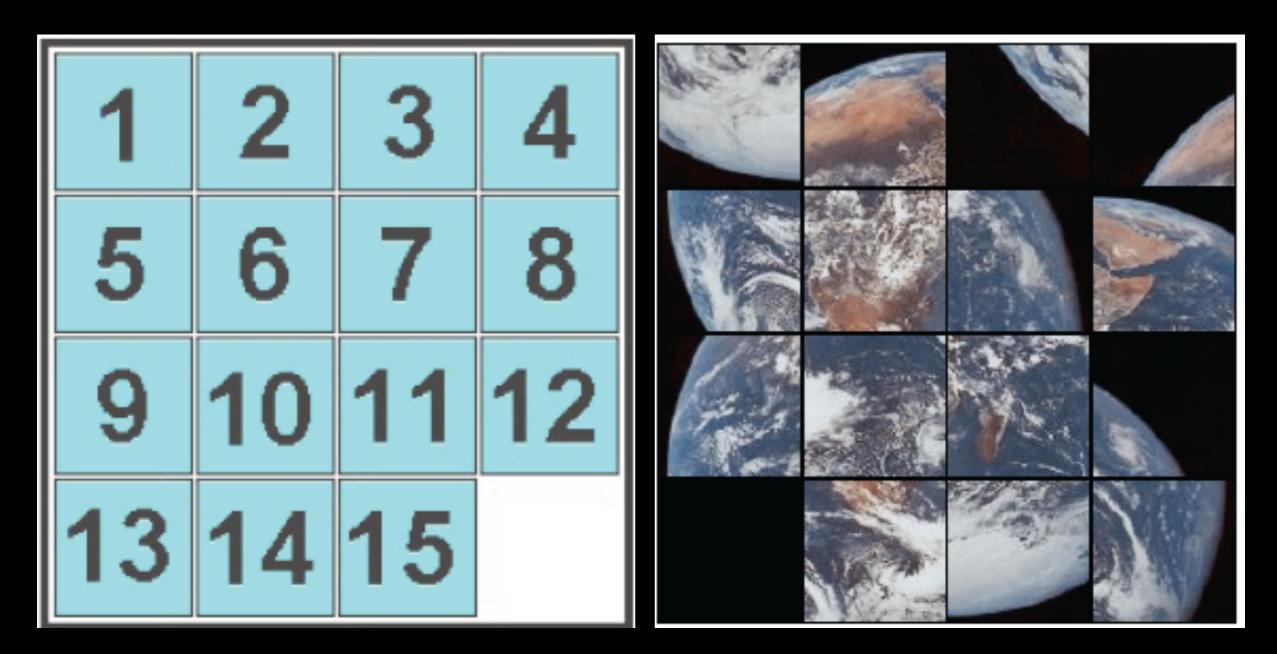
- through the crystal matrix (Nabarro-Herring)
- along the grain boundaries (Coble)







Diffusion creep

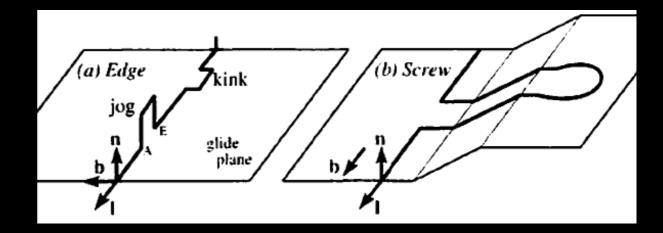




Dislocation creep

line defects move by dislocation

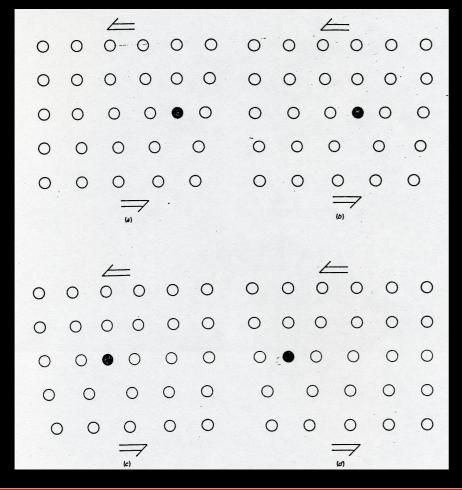
- two types of dislocations: "edge" and "screw"
- any line defect can be represented by linear combination of the two (simply add Burgers vectors)
- line dislocations have two types of motion: glide and climb
- independent of grain-size (point of difference with diffusion creep)

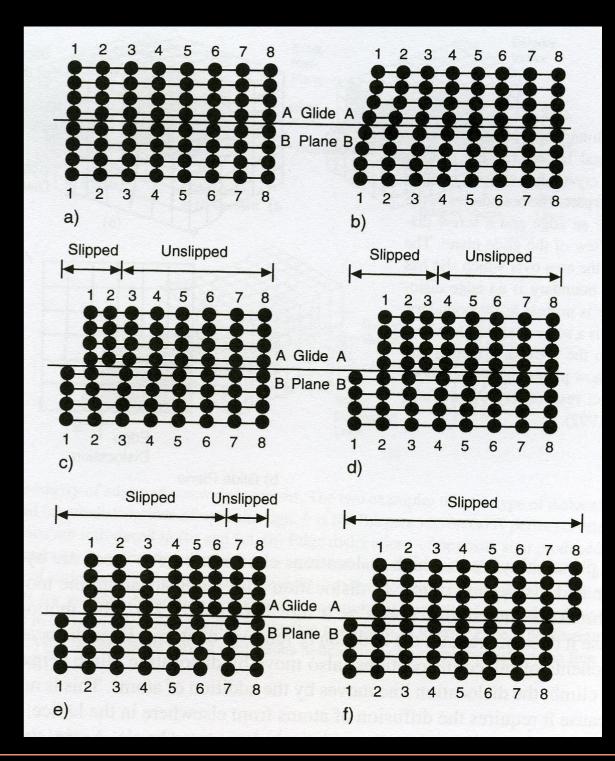




Glide process of dislocation creep

glide motion stays within glide plane

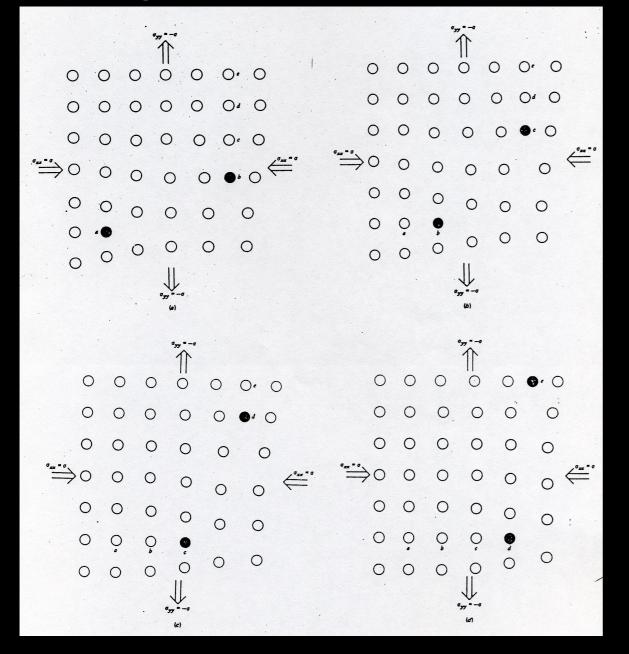


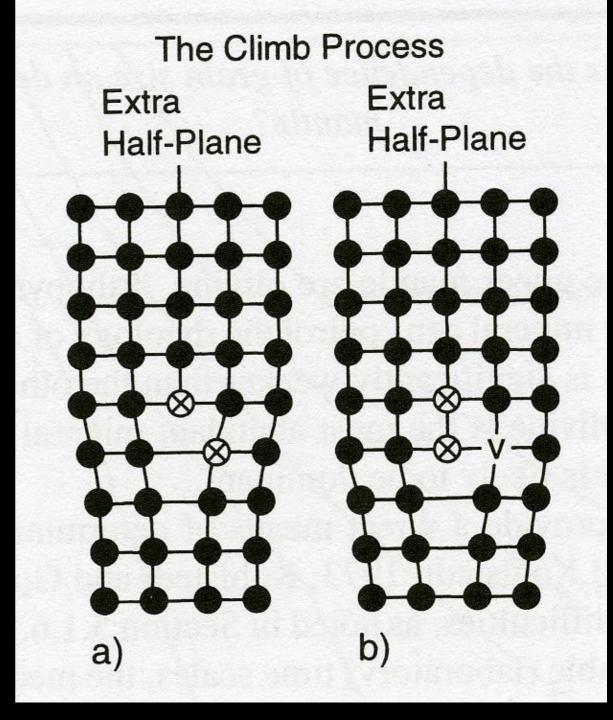




Climb process of dislocation creep

 climb dislocation occurs outside the glide surface





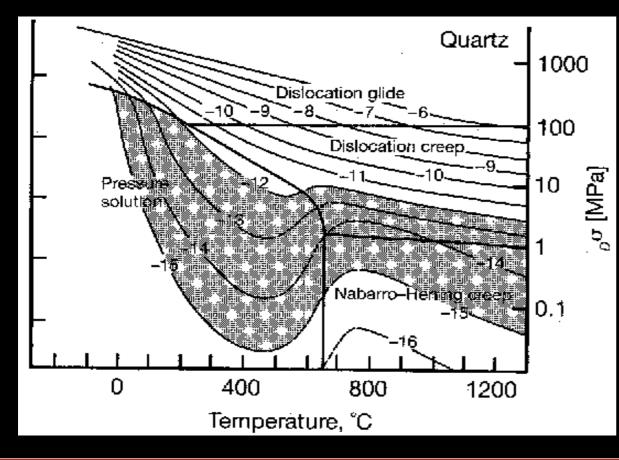


Deformation maps

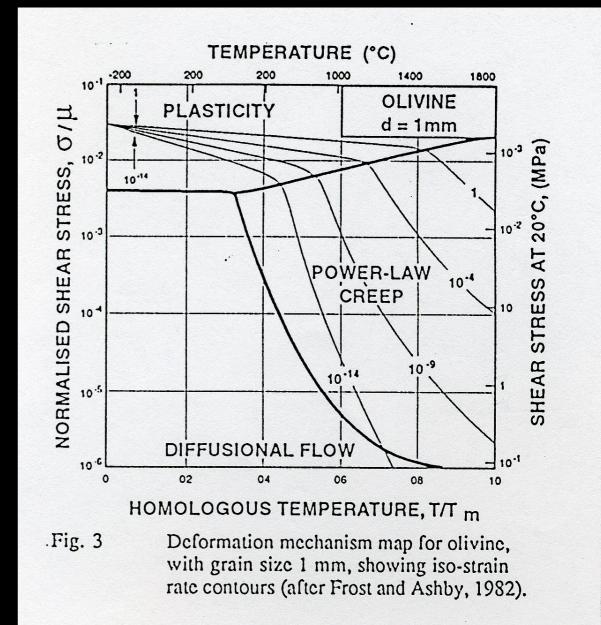
 for any given stress + temperature, one mechanism will be weaker (and preferred) over all others

$$\dot{\varepsilon} = \dot{\varepsilon}_{\text{disl}} + \dot{\varepsilon}_{\text{diff}}$$

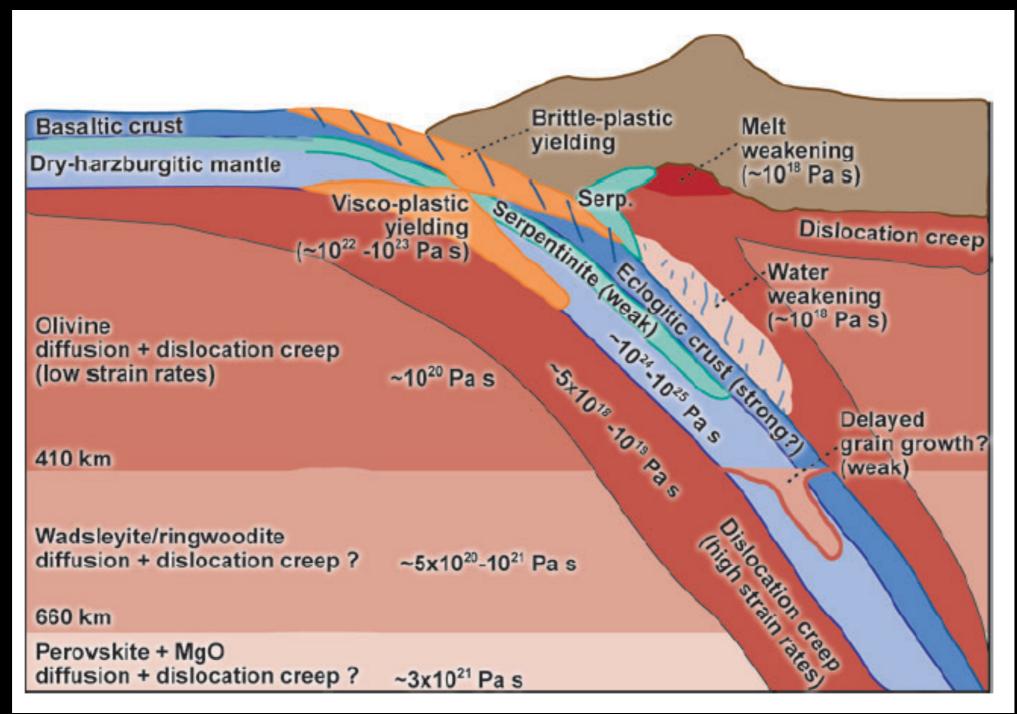
map assumes constant grain size



E TECTONIC



Creep mechanisms in the mantle



Billen, Annual Rev. Geophys., 2008



Flow rule

- relates the deformation (strain rate) to the applied deviatoric stress through a viscosity
- deformations add in series (viscosities add in parallel)

$\sigma = \eta \dot{\varepsilon}$	
$\dot{\varepsilon} = \frac{1}{\eta}\sigma$	
$\dot{\varepsilon}_{tot} = \dot{\varepsilon}_1$	$+\dot{\varepsilon}_2+\ldots$
1 = 1	1
$\eta_{eff} = \eta_1$	η_2 η_2



Flow rule

- for isotropic fluids, we can describe the flow rules with an effective strain rate, an effective stress, and an effective viscosity
- 2nd invariants of deviatoric stress and strain rates are scalars
- in practice, we know the strain-rates from the velocity field rather than stress, so viscosity is normally rewritten in terms of strain-rate

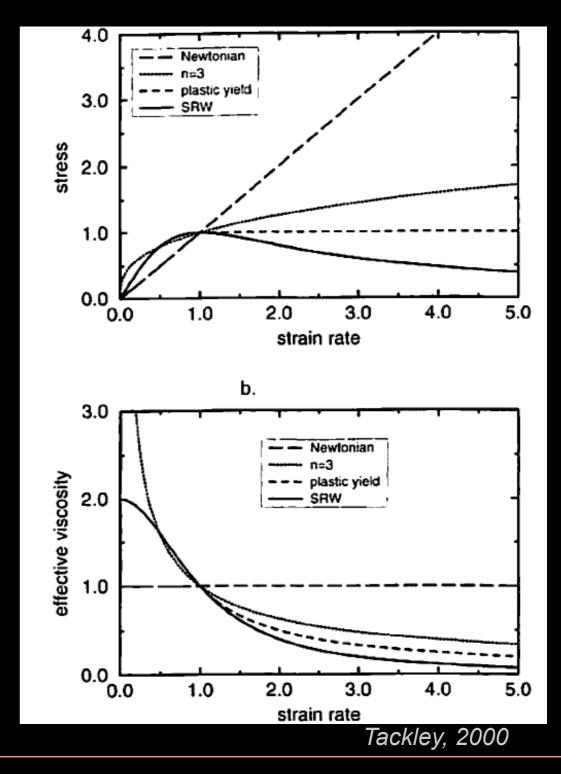
$$\sigma_{II} = \sqrt{\sigma} : \sigma$$
$$\dot{\varepsilon}_{II} = \sqrt{\dot{\varepsilon}} : \dot{\varepsilon}$$
$$\sigma_{II} = \eta_{eff} \dot{\varepsilon}_{II}$$

NOTE: contractions of these tensors usually have a 1/2 term times the sum of the squares of the components (when assuming co-axial compression / pure shear)



Rheology: land of jargon and confusion

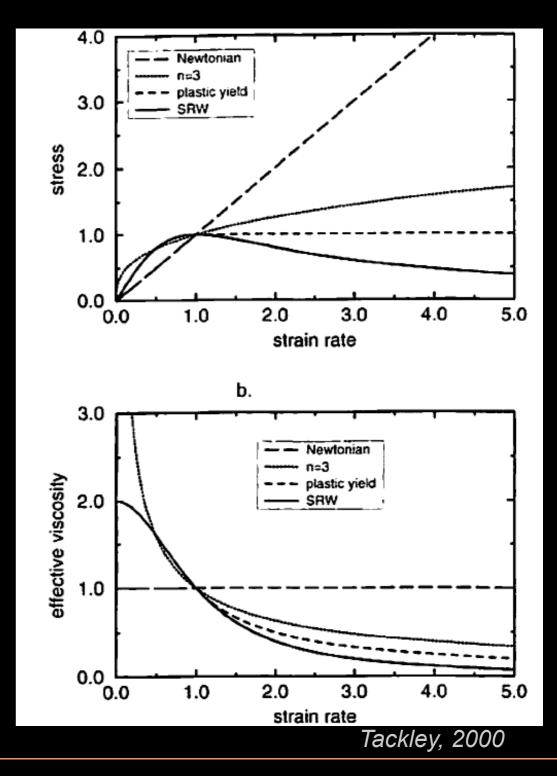
- Newtonian strictly means linear viscosity, but is commonly used to refer to diffusion creep with stress exponent (n=1)
- Non-Newtonian usually refers to non-linear rheology resulting from dislocation creep that is stress dependent through power law on stress with an exponent (n=3.5, or historically n=3)
- Example: diffusion creep is thermally activated and pressure dependent, and thus has exponential sensitivity to T,P but is sometimes referred to as "Newtonian" even though it is a non-linear f(T,P)





Comparison of deformation processes

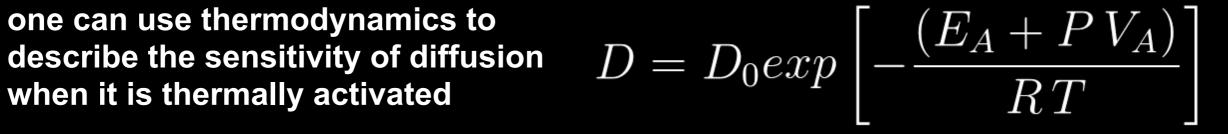
- any non-linear rheology can be represented by an effective linear rheology
- the effective viscosity is simply the slope of the line drawn from the origin to the stress-strainrate curve
- example of effective viscosity for several different rheologies that may be important for the mantle + plates
- think a little more about these curves stress vs. effort, efficiency vs. effort, and total productivity vs. effort





Arrhenius dependence

- one can use thermodynamics to
- the diffusion of vacancies, etc has an exponential dependence on T,P
- this can also be understood in terms of the homologous temperature
- the deformation resulting from the diffusion creep is the strain rate and is inversely proportional to viscosity
- the Arrhenius term describes the exponential behavior of viscosity
- also valid for dislocation creep



$$D = D_0 exp \left[-g \frac{T_m(P)}{T} \right]$$
$$\dot{\varepsilon} = \frac{1}{\eta} \sigma$$

$$\eta(T,P) = \eta_0 exp \left[\frac{(E_A + P V_A)}{R T} \right]$$



$$\dot{\varepsilon} = A\left(\frac{b^*}{d}\right)^m \left(\frac{\sigma}{\mu}\right)^n exp\left[-\frac{(E_A + P V_A)}{R T}\right]$$



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$$\dot{\varepsilon} = \frac{A}{\mu} \left(\frac{b^*}{d}\right)^m \sigma \left(\frac{\sigma}{\mu}\right)^{n-1} exp\left[-\frac{(E_A + PV_A)}{RT}\right]$$



$$\dot{\varepsilon} = A\left(\frac{b^*}{d}\right)^m \left(\frac{\sigma}{\mu}\right) \left(\frac{\sigma}{\mu}\right)^{n-1} exp\left[-\frac{(E_A + PV_A)}{RT}\right]$$
$$\frac{1}{\eta} = \frac{A}{\mu} \left(\frac{b^*}{d}\right)^m \left(\frac{\sigma}{\mu}\right)^{n-1} exp\left[-\frac{(E_A + PV_A)}{RT}\right]$$



$$\frac{1}{\eta} = \frac{A}{\mu} \left(\frac{b^*}{d}\right)^m \left(\frac{\sigma}{\mu}\right)^{n-1} exp\left[-\frac{(E_A + P V_A)}{R T}\right]$$
$$\eta = \frac{\mu}{A} \left(\frac{d}{b^*}\right)^m \left(\frac{\mu}{\sigma}\right)^{n-1} exp\left[\frac{(E_A + P V_A)}{R T}\right]$$

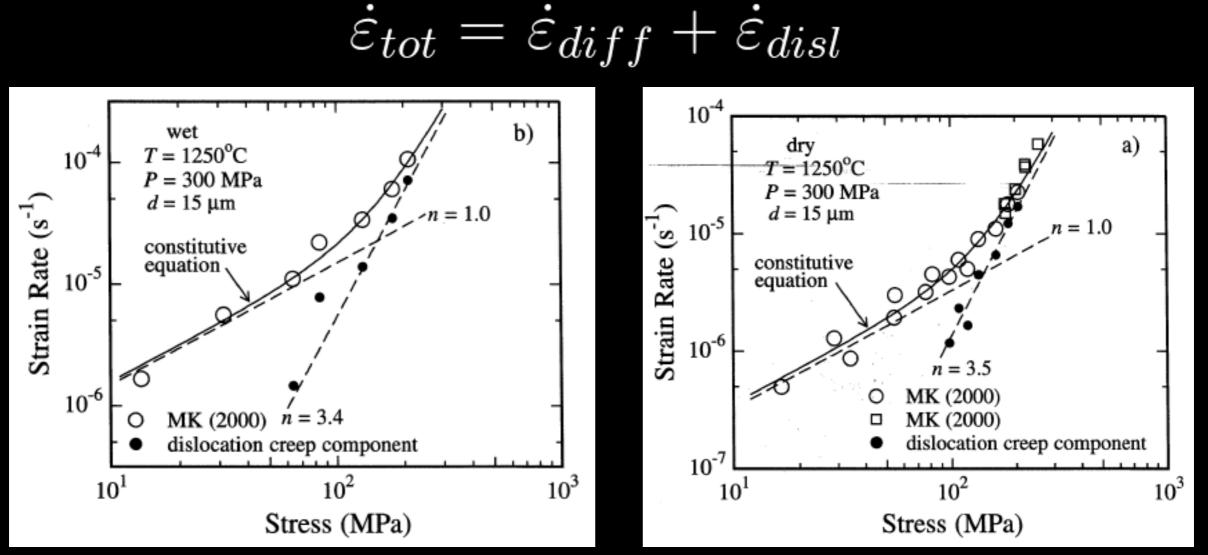


$$\eta = \frac{\mu}{A} \left(\frac{d}{b^*}\right)^m \left(\frac{\mu}{\sigma}\right)^{n-1} exp\left[\frac{(E_A + PV_A)}{RT}\right]$$
$$\dot{\varepsilon}_{tot} = \dot{\varepsilon}_{diff} + \dot{\varepsilon}_{disl}$$
$$\frac{1}{\eta_{eff}} = \frac{1}{\eta_{diff}} + \frac{1}{\eta_{disl}}$$

note: diffusion and dislocation creep have different activation enthalpies



What is "n" for dislocation creep?

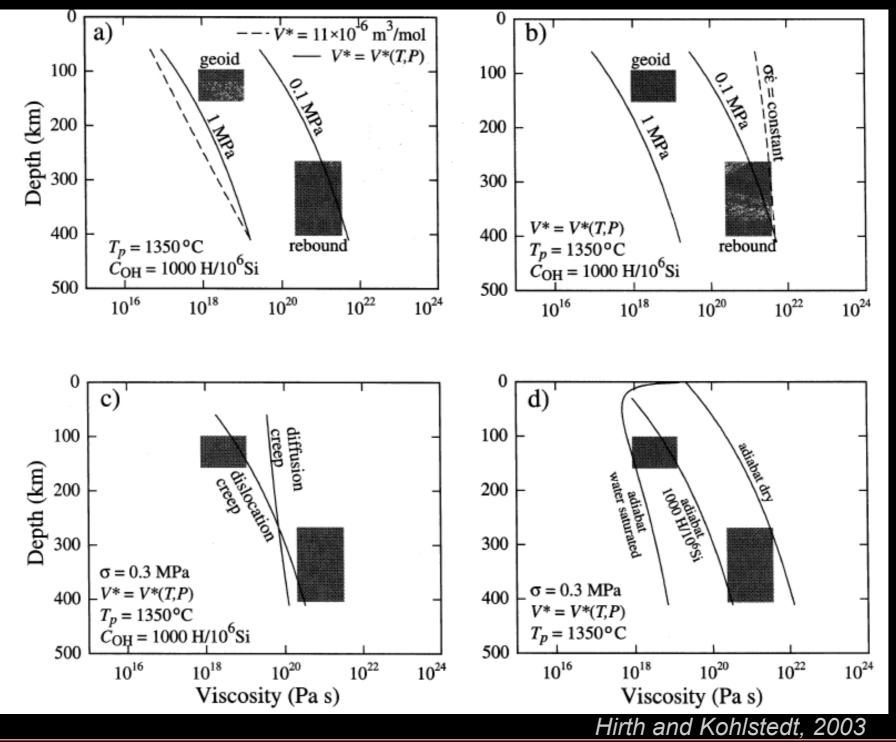


Hirth and Kohlstedt, 2003

 note: different y-intercept corresponds to different activation enthalpies which is due to the influence of water



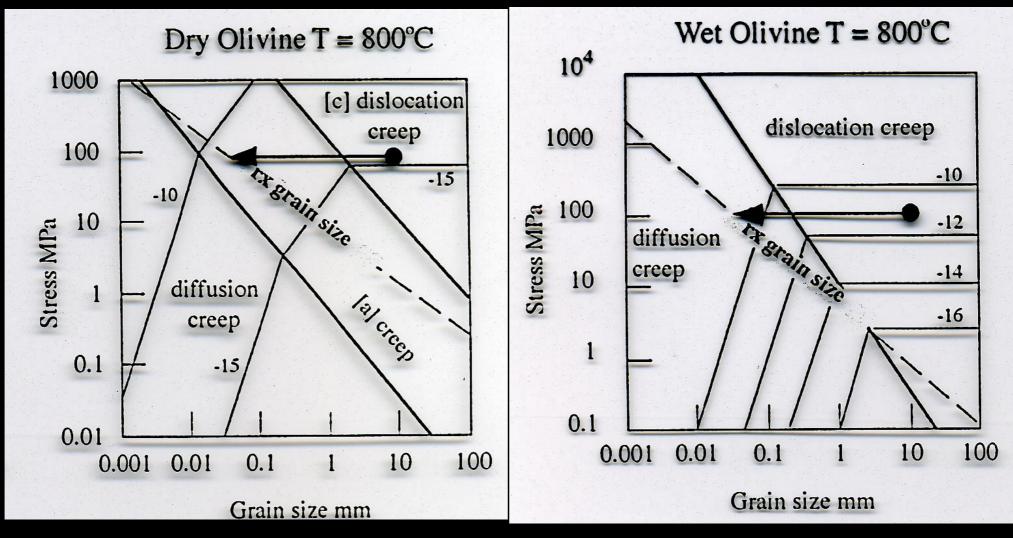
Extrapolating to mantle conditions





Just add water...

- contour lines are strain rates in (1/s) dislocation is horizontal because it is independent on grain size but diffusion creep shows grain size dependence
- much lower stresses required to achieve same strain rate when water is
 present (note different y-axis) which translates into lower activation enthalpy





Just add water...

- contour lines are strain rates in (1/s) dislocation is horizontal because it is independent on grain size but diffusion creep shows grain size dependence
- much lower stresses required to achieve same strain rate when water is present (note different y-axis)
- translates into lower activation energy and volume (E* and V*)

$$\dot{\varepsilon}(C_{OH}) = A\left(\frac{b^*}{d}\right)^m \left(\frac{\sigma}{\mu}\right)^n exp\left[-\frac{(E_A^* + P V_A^*)}{R T}\right]$$

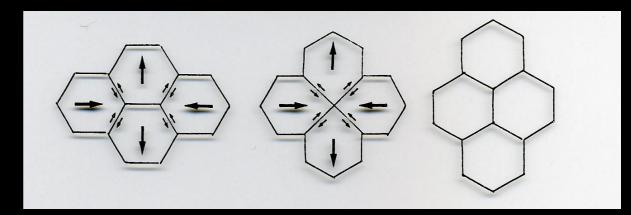


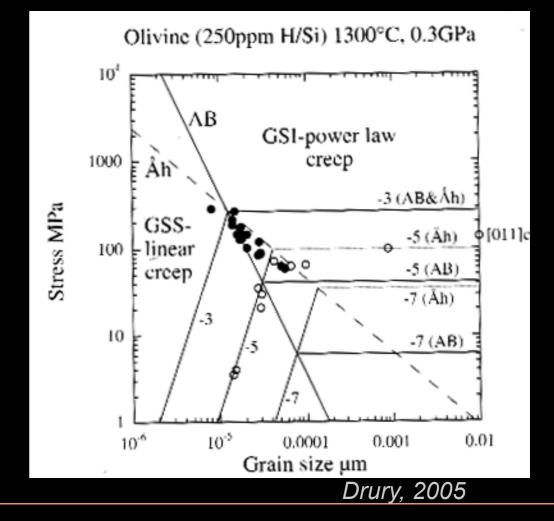
Grain size sensitive creep

 GSS creep relies on combination of mechanisms: grain boundary sliding (GBS) and basal slip along the easiest slip system (easy)

$$\dot{\varepsilon} = \dot{\varepsilon}_{\text{disl}} + \dot{\varepsilon}_{\text{diff}} + [1/\dot{\varepsilon}_{\text{gbs}} + 1/\dot{\varepsilon}_{\text{easy}}]^{-1}$$

- GSS operates at small grain size
 only relevant to diffusion creep
- accommodates large amounts of strain without deforming crystals (referred to as superplasticity)
- GSS important process in lower mantle with small grain size (Karato et al., Science, 1995) through extrapolation to high P, low strainrate



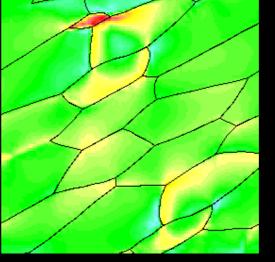


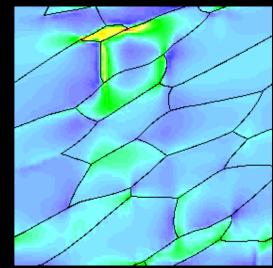


Recrystallization

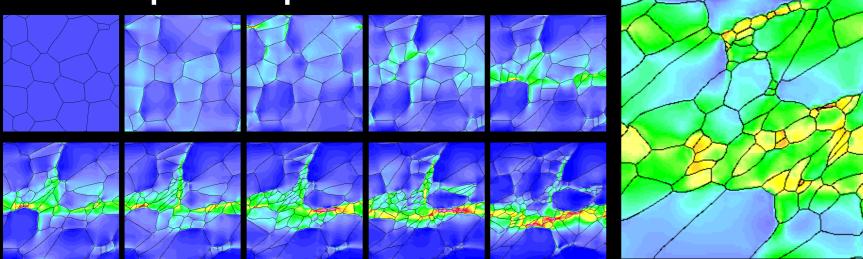
- dynamic process changing the distribution of grain sizes
- grain size reduction (large grains fracture) and growth (small grains fuse together)
- microstructure important feedback with dislocation creep
- time-dependent process

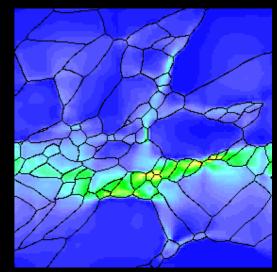
linear (n=1) power law (n=3.5) without recrystallization





with recrystallization



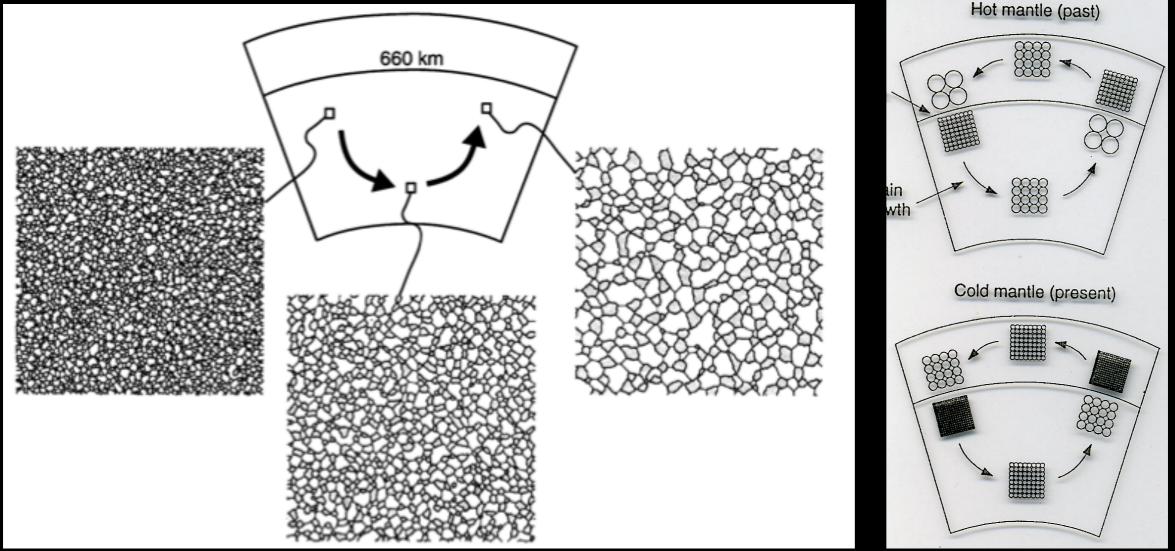


Jessell et al, EPSL, 2005



Grain growth in the lower mantle

- recrystallization for material going down through 660 phase change
- diffusion creep likely dominant mechanism for lower mantle but if the grain size is sufficiently small, then superplasticity could be important mechanism



Solomatov, EPSL, 2001



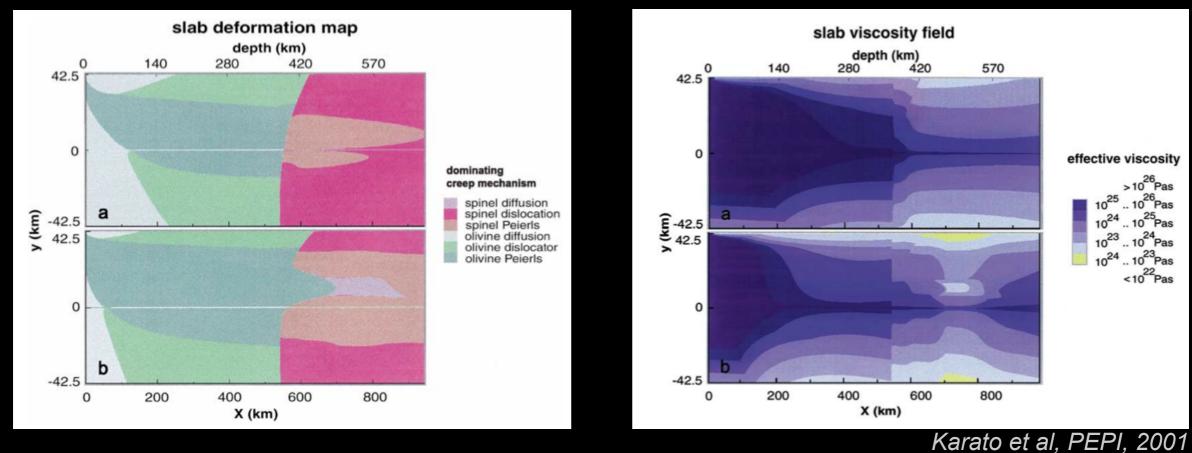
Harper-Dorn creep

low-temperature plasticity OR Peierls stress mechanism*

$$\dot{\varepsilon} = A\left(\frac{b^*}{d}\right)^m \left(\frac{\sigma}{\mu}\right)^n exp\left[-\frac{(E_A + P V_A)}{R T} \left(1 - \frac{\sigma}{\sigma_p}\right)^q\right]$$

*read about Peierls other contributions to science here: https://en.wikipedia.org/wiki/Rudolf_Peierls

applicable to the interior of a subducted slab (below: 4 cm/yr and 10 cm/yr)





Generalized flow rule for a slab

$$\frac{\dot{\varepsilon}_{tot} = \dot{\varepsilon}_{diff} + \dot{\varepsilon}_{disl} + \dot{\varepsilon}_{H-D}}{\eta_{eff}} = \frac{1}{\eta_{diff}} + \frac{1}{\eta_{disl}} + \frac{1}{\eta_{H-D}}$$

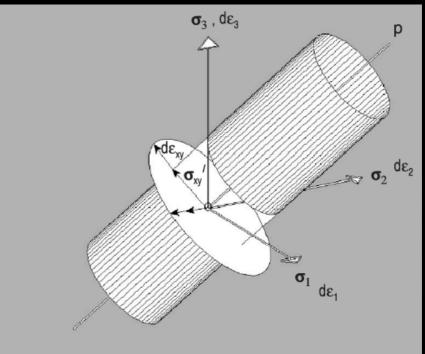
- diffusion, dislocation, Harper-Dorn creep mechanisms work independently, so in general, only a single power-law (m,n,q) is in effect at any given time
- usual values for exponents (m,n,q) are (2.5,1,0) + (0,3.5,0) + (0,2,2)
- each mechanism has a different value for A; E and V depend on water content

$$\dot{\varepsilon} = A\left(\frac{b^*}{d}\right)^m \left(\frac{\sigma}{\mu}\right)^n exp\left[-\frac{(E_A + P V_A)}{R T} \left(1 - \frac{\sigma}{\sigma_p}\right)^q\right]$$



Viscoplasticity

- most materials have finite strength described by limiting stress and materials cannot support stresses in excess of their yield stress
- upon reaching their yield stress they deform through plastic flow (a solid beam starts to act like toothpaste)

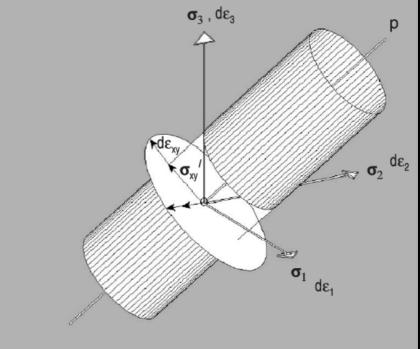


von Mises yield envelope



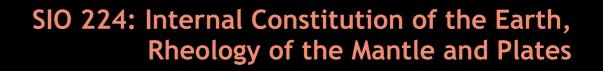
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- actually, care must be taken to guarantee one is actually on the yield surface and remains on it (i.e. use harmonic avg at your own risk)



von Mises yield envelope

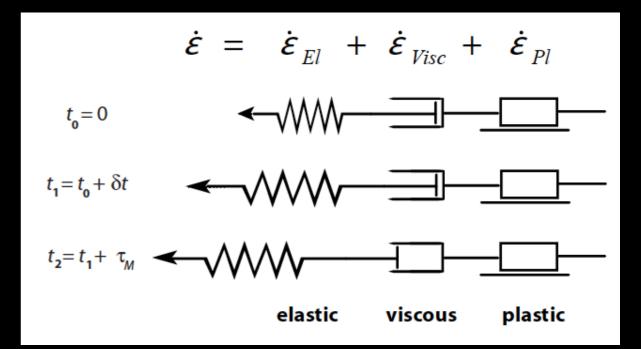
$$\eta = \begin{cases} \eta = \frac{\sigma_{II}}{\dot{\varepsilon}_{II}} & (\sigma < \sigma_{yield}) \\ \eta_{\text{eff}} = \frac{\sigma_{yield}}{\dot{\varepsilon}_{II}} & (\sigma \ge \sigma_{yield}) \end{cases}$$





Viscoplasticity

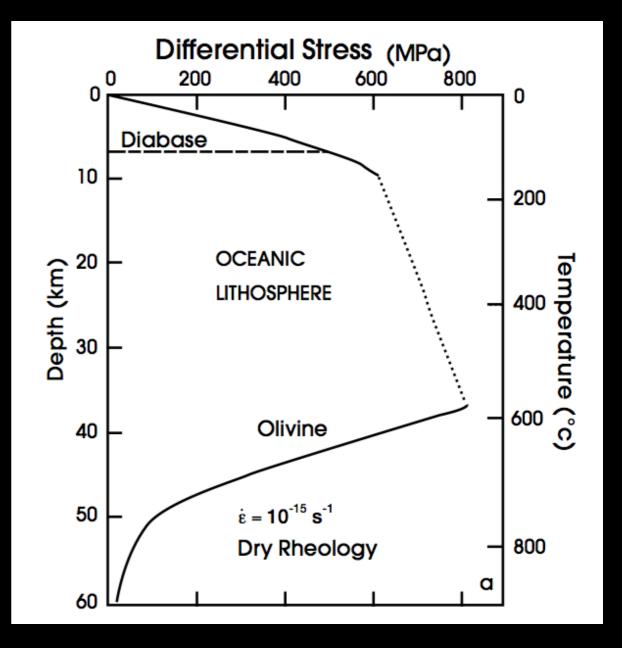
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- upon reaching their yield stress they deform through plastic flow (a solid beam starts to act like toothpaste)
- actually, care must be taken to guarantee one is actually on the yield surface and remains on it (i.e. use harmonic avg at your own risk)
- visco-elasto-plastic behavior can be written as generalized flow rule with an associated flow for each of the viscous, elastic, and plastic parts



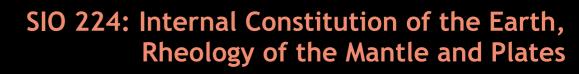


Strength envelope for an oceanic plate

- maximum stress that is supported is equivalent to a yield stress
- three layer lithosphere: brittle crust, strong core, ductile underside
- based on extrapolating the deformation behavior of crystals (microscale) to that of a rock scale (macroscale)
- role of large scale faults and tectonic fabric (mesoscale)
- likely mechanism in strong core is Peierls creep (low strainrate, high stress)



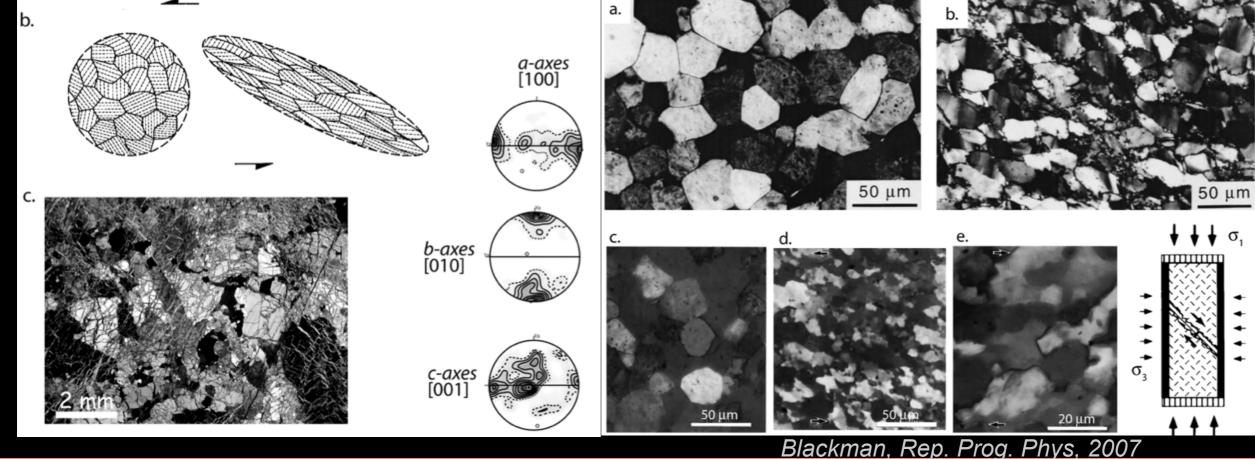
Kohlstedt et al, JGR, 1995





Anisotropy

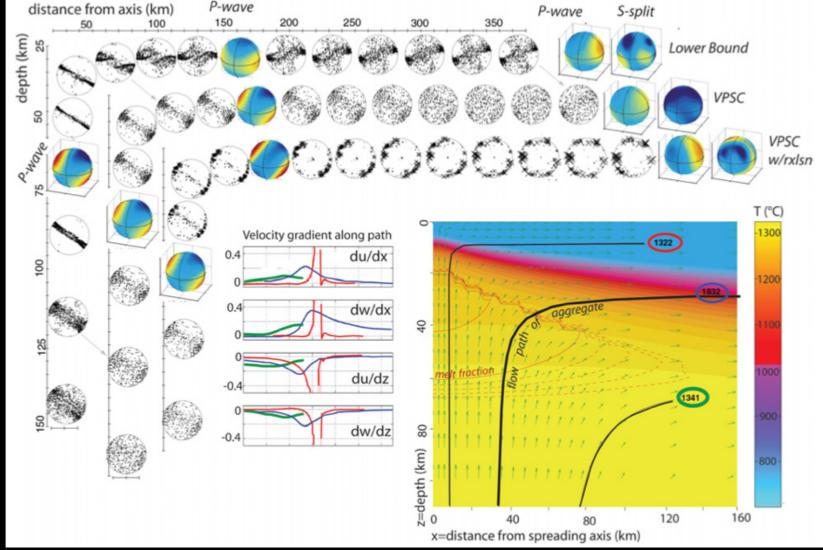
- dislocation creep (n=3.5) and dynamic recrystallization during deformation generate crystal alignments and lattice preferred orientation (LPO) of crystals
- significant seismic anisotropy observed in the upper mantle -> implies dislocation creep is the dominant mechanism
- lack of anisotropy seen in lower mantle -> diffusion creep is dominant





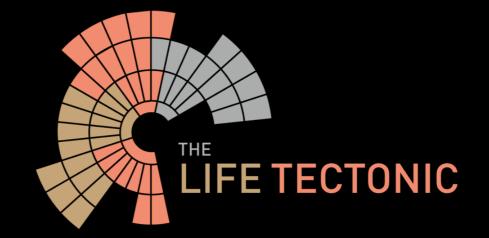
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Blackman, Rep. Prog. Phys, 2007





www.thelifetectonic.com



Thank you! Questions??