

# Identifying landslides from continuous seismic surface waves: a case study of multiple small-scale landslides triggered by Typhoon Talas, 2011

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## SUMMARY

Landslides can cause devastating damage. In particular, heavy rainfall-triggered landslides pose a chain of natural hazards. However, such events are often difficult to detect, leaving the physical processes poorly understood. Here we apply a novel surface-wave detector to detect and locate landslides during the transit of Typhoon Talas 2011. We identify multiple landslides triggered by Typhoon Talas, including a landslide in the Tenryu Ward, Shizuoka prefecture, Japan,  $\sim 400$  km east from the typhoon track. The Tenryu landslide displaced a total volume of  $1.2 - 1.5 \times 10^6$  m<sup>3</sup>. The landslide is much smaller than those detected by using globally recorded surface waves, yet the event generated coherent seismic signals propagating up to 3000 km away. Our observations show that attributes of small and large landslides may follow the same empirical scaling relationships, indicating possible invariant failure mechanisms. Our results also suggest an alerting technology to detect and locate landslides with a sparse seismic network.

**Key words:** Time-series analysis; Seismic noise; Surface waves and free oscillations; Wave propagation; Tectonics and landscape evolution.

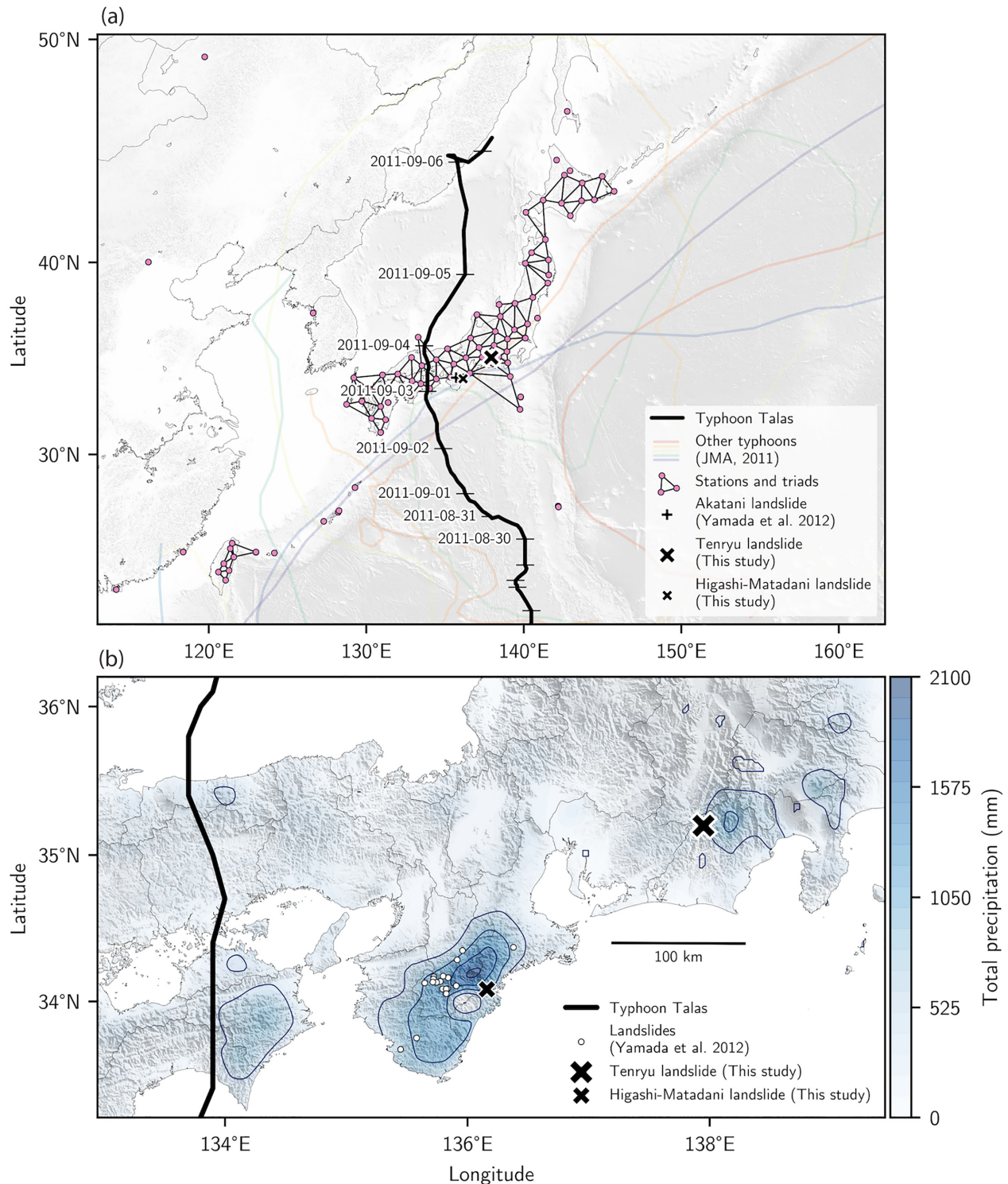
## 1 INTRODUCTION

Deep-seated catastrophic landslides can displace mass over a large range of volumes rapidly and can cause significant hazards to mountain communities and infrastructure (Spiker & Gori 2003; Hewitt *et al.* 2008; Hibert *et al.* 2011; Chigira *et al.* 2013; Ekström & Stark 2013). Mitigations of such disastrous events rely on robust monitoring of landslide failure processes, yet most observations of landslide dynamics remain retrodictive. Broadband seismic observations can help detecting and locating these events even when landslides are distant from the seismic networks (Ekström & Stark 2013; Fan *et al.* 2020).

Landslides can generate broad-band seismic signals (Kanamori & Given 1982; Kawakatsu 1989; Brodsky *et al.* 2003; Allstadt 2013; Hibert *et al.* 2015). Short-periods ( $< 1$  s) (Hibert *et al.* 2011; Yamada *et al.* 2012; Doi & Maeda 2020) and intermediate- to long-periods (20–150 s, Moretti *et al.* 2012; Allstadt 2013; Ekström & Stark 2013; Li *et al.* 2019; Zhang *et al.* 2019) seismic signals are commonly used for detecting landslides and studying landslide dynamics. For example, short-period signals have proven efficient for detecting and evaluating landslides (Dammeier *et al.* 2016; Manconi *et al.* 2016; Chao *et al.* 2017; Dietze *et al.* 2017; Fuchs

*et al.* 2018). However, such operations are often limited to local or regional distances due to seismic attenuation. The intermediate- to long-period (35–150 s) seismic surface waves are the primary means to detect and locate distant landslides (Ekström 2006; Ekström & Stark 2013). For example, Rayleigh waves have proven effective for detecting teleseismic landslides (Ekström 2006; Lin *et al.* 2010). These landslides can displace  $\geq 2 \times 10^{10}$  kg rocks and generate surface waves with amplitudes equivalent to those of  $M \geq 4.6$  earthquakes (Ekström 2006; Ekström & Stark 2013). In contrast, smaller size landslides are infrequently reported from surface wave detectors, leaving their occurrence poorly understood. Recently, automatic classifiers show promises in detecting small landslides from continuous regional seismic records (e.g. Dammeier *et al.* 2016; Hibert *et al.* 2017; Provost *et al.* 2017). However, such algorithms have not been applied to systematically locate landslides due to the limited seismic network coverage (Hibert *et al.* 2019).

The 2011 Typhoon Talas brought precipitation exceeding 2000 mm and caused 50+ landslides adjacent to the typhoon track in Nara, Wakayama and Mie prefectures in western Japan (Yamada *et al.* 2012; Chigira *et al.* 2013) (Fig. 1). Among them, 18 landslides were detected and located by using short-period



**Figure 1.** Overview of the study area. (a) Map shows the available seismic stations during the study period, the track of Typhoon Talas and the landslide locations. Background topography/bathymetry are from the GEBCO 2019 Grid (GEBCO Bathymetric Compilation Group 2019 2019). (b) Background colour is the total precipitation during 30 August 2011 to 6 September 2011 observed at the Automated Meteorological Data Acquisition System (AMeDAS) stations. The blue contour denotes every 500 mm total precipitation. The grey lines denote the administrative boundaries.

(0.25–1 s) seismic records near the landslide sources (Yamada *et al.* 2012). However, due to seismic attenuation of these waves, such procedure is inadequate to detect landslides that were away from the Typhoon track. Intriguingly, the precipitation in Shizuoka

prefecture is over 1000 mm, which is  $\sim 400$  km away from the typhoon track (Fig. 1b). However, no landslides were reported in this region by previous seismic studies (e.g. Yamada *et al.* 2012).

Here we apply a surface wave detector that is based on the AELUMA method (Automated Event Location Using a Mesh of Arrays, de Groot-Hedlin & Hedlin 2015; Fan *et al.* 2018) to investigate landslide activities across Japan during the transit of Typhoon Talas. This method has been applied to the USArray with over 400 stations and located various unconventional seismic sources (Fan *et al.* 2018, 2019, 2020). In this study, we identify three new landslides, including one in Tenryu, Shizuoka prefecture, which is 400 km away from the track of Typhoon Talas. The landslide generates coherent surface wavefields that are recorded by stations across Japan and Taiwan but only displace a total volume of  $1.2\text{--}1.5 \times 10^6 \text{ m}^3$  (Kanto Regional Forest Office Japan 2012; Seo *et al.* 2012; Yumoto & Takashima 2013). The results show promises of future near-real-time monitoring of landslide activities in Japan.

## 2 DATA

We use continuous seismic data from 103 stations of the National Research Institute for Earth Science and Disaster Resilience F-net (NIED 2019) and the Broadband Array in Taiwan for Seismology TW (IES 1996) networks shown in Fig. 1(a). We download the vertical-component long-period (1-s-sampled LHZ) records of 3 to 4 September 2011, during Typhoon Talas' transit in Japan (Fig. 1a, Yamada *et al.* 2012). We then remove the instrumental response to utilize data from different instruments. The records are bandpass filtered at 20–50 s with a 4th-order non-causal Butterworth filter.

## 3 METHOD

### 3.1 Detecting and locating seismic sources using seismic surface waves

We apply the AELUMA-based surface wave detector to detect and locate seismic events. Following de Groot-Hedlin & Hedlin (2015), we first divide the 103 stations into non-overlapping 68 triangular subarrays (triads), and remove triads with internal angles beyond the range of  $30^\circ$  to  $120^\circ$  (Fig. 1a, Lee & Schachter 1980; Thompson & Shure 2016). Secondly, we apply tau-p beamforming analysis for each triad, and measure relative traveltimes between station pairs of coherent signals to solve for a centroid arrival time and a propagation direction (Fig. S1). Thirdly, the detections are aggregated into non-overlapping clusters. We then invert the seismic source locations with the clusters by grid-searching possible source locations (Fan *et al.* 2018). To neutralize off-great-circle path propagation effects, we also apply empirical calibrations from measurements of earthquakes in the Global Centroid Moment Tensor (GCMT) project (Dziewonski *et al.* 1981; Ekström *et al.* 2012) and landslides reported in a previous study (Yamada *et al.* 2012). After obtaining the source locations, we finally perform a quality control step to discard sources detected by less than 10 triads. These empirical parameters are different than those applied to the USArray (e.g. Fan *et al.* 2018), but comparable parameters were examined in de Groot-Hedlin & Hedlin (2018) and proven effective. Details of the algorithm are described in Fan *et al.* (2018) and de Groot-Hedlin & Hedlin (2015).

### 3.2 Centroid-single force modelling

To investigate the source mechanisms of the newly identified seismic sources (e.g. seismic event E1, Fig. 2, Table 1), we perform the seismic waveform inversion to model the source as centroid-single

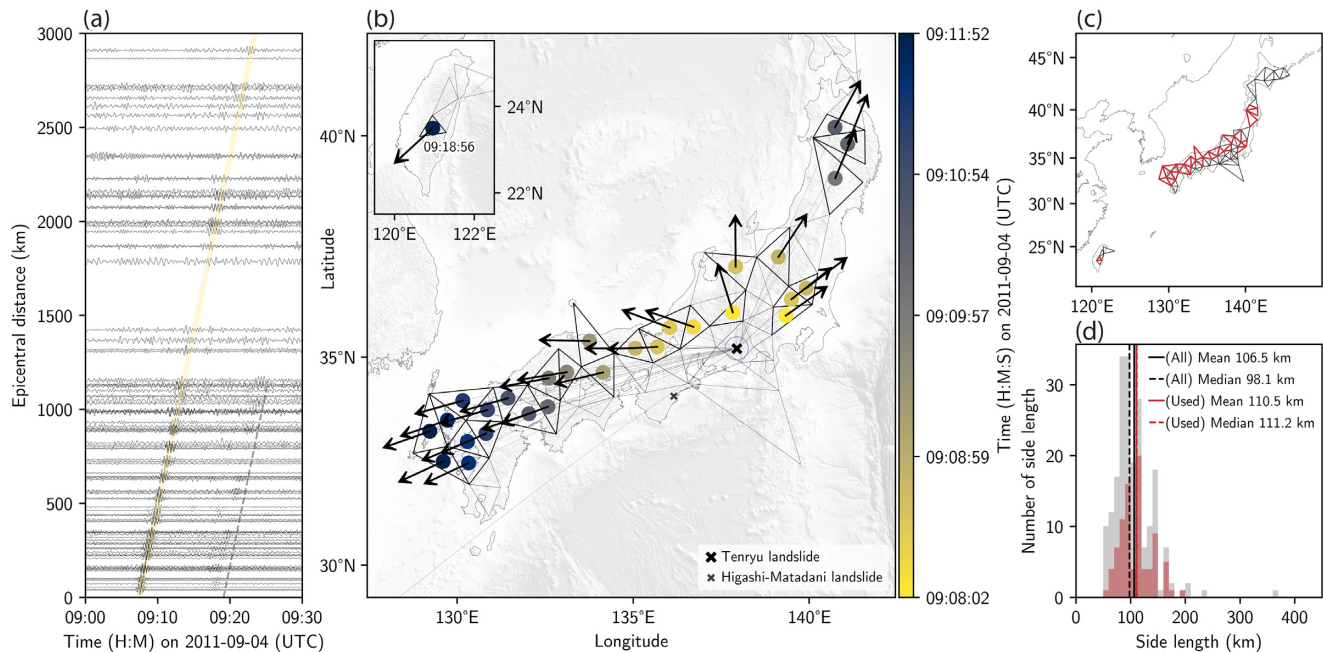
forces (CSF, Kawakatsu 1989; Tsai & Ekström 2007; Ekström & Stark 2013). As discussed later, our newly identified seismic sources (e.g. seismic event E1, Fig. 2, Table 1) are likely landslides, which show clear seismic surface waves in a narrow intermediate period band (20–50 s) but do not show clear *P*- or *S*-arrivals (Figs 3, 4 and S5). We adopt a conventional time-domain method (e.g. Fan *et al.* 2020) to obtain a CSF model of the seismic source. The method assumes the force functions at three directions are equal-time length symmetric boxcar functions, representing a constant initial acceleration and an equal-duration, equal-amplitude arresting deceleration. For a candidate CSF model, we calculate the associated synthetic seismograms by convolving the model with Green's functions, and evaluate the model in the time domain for an average minimum  $\ell_2$  waveform-misfit (Fan *et al.* 2020). We construct CSF models to explain both the Rayleigh and Love waves (20–50 s) at stations within 150 km (e.g. Fig. 5a). In practice, we grid-search the source duration and the three centroid force amplitudes in a 4-D model-space with ranges from  $-2 \times 10^{10}$  to  $2 \times 10^{10}$  N for the up–down, north–south and east–west components, and from 10 to 50 s for the duration. The search steps are  $0.05 \times 10^{10}$  N for up–down and east–west components and  $0.005 \times 10^{10}$  N for a north–south component. We test durations of 10, 20, 24, 30, 34, 40 and 50 s. We compute nine-components of force-source Green's functions for each source–station pair at the three directions using the Instaseis method (van Driel *et al.* 2015). The Instaseis method uses a pre-computed Green's function database, which is calculated by the axisymmetric spectral-element code AxisEM with the anisotropic version of the PREM model up to 5 s (Dziewonski & Anderson 1981; Nissen-Meyer *et al.* 2014). This allows efficient computation of broadband synthetic seismograms and a relative complete search of the model space. This Green's function database can be directly obtained from IRIS DMC Syngine (Krischer *et al.* 2017). The regional geological structures are complex and can have strong velocity effects. Such effects dominantly show as arrival time anomalies. We correct the crust-heterogeneity effects by cross-correlating the observed and synthetic waveforms and aligning them before computing the waveform misfit. Both the observed and synthetic waveforms are resampled at 1 s and filtered at 20–50 s with a 4th-order Butterworth bandpass filter before the waveform misfit calculation. Even though source kinematics can be rather complex (Moretti *et al.* 2012, 2020; Yamada *et al.* 2013, 2018), the simple boxcar landslide models are representative of the ground loading and unloading processes for the sliding processes because of us focusing on the intermediate-period (20–50 s) surface waves. As shown in the sections below, the models prove sufficient in explaining the seismic observations well (Fig. 5c).

## 4 RESULTS

### 4.1 Overview of the detected seismic events

We initially located 25 seismic events from 3 to 4 September 2011 (Table S1). We further screen the events by visually inspecting the waveform records aligned with the source epicentres, and events that generate coherent wave trains are kept for further analysis. This quality control step removes nine events (e.g. Fig. S2) and leaves 16 candidate events for further analysis (e.g. Fig. 2a). Thirteen of the candidate events are earthquakes in standard earthquake catalogues (Dziewonski *et al.* 1981; Japan Meteorological Agency 2011; Ekström *et al.* 2012; U.S. Geological Survey Earthquake Hazards Program 2017) and two sources were landslides reported in Yamada





**Figure 2.** Detection and location of the Tenryu landslide. (a) Self-normalized bandpass-filtered (20–50 s) waveforms aligned by the epicentre of the Tenryu landslide (E1). The yellow line shows the reference wavefront travelling at a phase velocity of  $3.11 \text{ km s}^{-1}$ . The dashed line indicates wave trains travelled from the Higashi-Matadani landslide (E2). (b) The thick and thin triangles are the triad subarrays. The arrow is the observed arrival angle. The colour for each dot represents the observed arrival time. The thin line between the epicentre and the centroid of each triad is the great circle path. The blue ellipse denotes the estimated location uncertainty. Inset is the triad measurement in Taiwan for the Tenryu landslide. (c) Black triangles are the triads available on 4 September 2011. Red triangles are the triads used for detection of the Tenryu landslide (E1). (d) Histogram of the side length of the triads every 10 km bin.

**Table 1.** List of landslides identified by this study.

Landslide name	Time (UTC)	Location	Reference
Tenryu (E1)	2011-09-04 09:07:28	35.1992°N, 137.9479°E	Newly identified by this study
Higashi-Matadani (E2)	2011-09-04 09:16:58	34.0823°N, 136.1602°E	Newly identified by this study
pre Higashi-Matadani (E3)	2011-09-04 09:16:55?	34.0823°N?, 136.1602°E?	Newly identified by this study
Ohto-Shimizu	2011-09-03 22:06:38	34.0447°N, 135.2156°E	This study and Yamada <i>et al.</i> (2012)
Akatani	2011-09-04 07:22:11	34.1557°N, 135.5472°E	This study and Yamada <i>et al.</i> (2012)

*et al.* (2012) (Table 1). We find one new unknown seismic event (E1) based on the initial set of parameters (Fig. 2b). As we detail in the later sections, we also observe weaker coherent phases following this new unknown source (Figs 2a and S4a). After re-examining the propagation direction and centroid time measurements, we identify two more events (E2 and E3) related to those signals, which are also absent in the standard earthquake or landslide catalogues (Dziewonowski *et al.* 1981; Japan Meteorological Agency 2011; Ekström *et al.* 2012; Yamada *et al.* 2012; U.S. Geological Survey Earthquake Hazards Program 2017). Thus in total, we identify three new seismic events (Table 1).

## 4.2 Major landslide E1

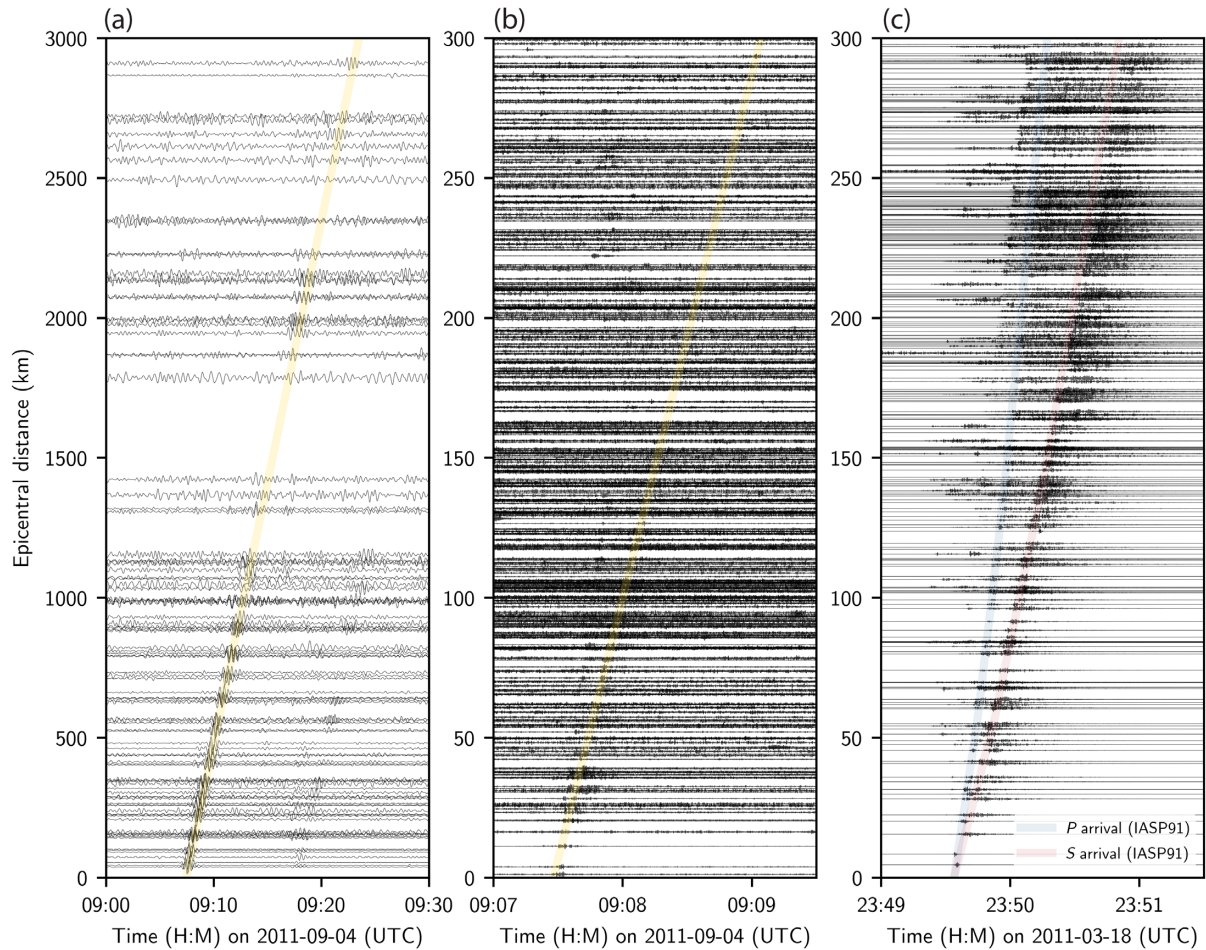
Seismic event E1 occurred on 4 September 2011, 09:07:28 (UTC) in Tenryu Ward, Shizuoka prefecture, Japan (35.1992°N, 137.9479°E, Fig. 2b). The waveform record-section of E1 shows a coherent wavefield propagating up to 3000 km with an estimated phase velocity of  $3.11 \text{ km s}^{-1}$  (Fig. 2a). The E1 location is resolved from measurements of 29 triads, including one in Taiwan (2000 km away from the epicentre) (Fig. 2b). The location uncertainty of E1 is  $\sim 30 \text{ km}$  (Fig. 2b), which is about one grid separation ( $\sim 30 \text{ km}$ ) (Fan *et al.* 2018). The surface-wave magnitude ( $M_{\text{SW}}$ , Ekström 2006) of the

event is 4.3. Our preferred CSF model of the E1 event has a misfit reduction of 72 per cent with peak force amplitudes of  $0.55 \times 10^{10} \text{ N}$ ,  $0.055 \times 10^{10} \text{ N}$  and  $0.6 \times 10^{10} \text{ N}$  for the up–down, north–south and east–west components, respectively (Fig. 5b). The preferred model suggests a source duration of 20 s and the sharp increase of the data misfit for models of longer durations suggests that the E1 event evolved rapidly (Table 2).

## 4.3 Minor landslide E2 and E3

From the record section in Fig. 2(a), we observe a weaker coherent phase  $\sim 10 \text{ min}$  after the E1 event. We re-examine the propagation direction and centroid time measurements and locate another event, E2, with only seven triads. Seismic event E2 is located near Higashi-Matadani in Mie prefecture (34.0823°N, 136.1602°E), occurring on 09:16:58 (UTC), 4 September 2011 (Fig. S4b) with a location uncertainty of  $\sim 30 \text{ km}$ . This event is adjacent to the Ohtaki landslide identified in Yamada *et al.* (2012) but occurred one hour later. There was a Japan Meteorological Agency (JMA) magnitude ( $M_{\text{JMA}}$ ) 1.7 earthquake in the area, but the near-field short-period records show that E2 was not the  $M_{\text{JMA}}$  1.7 earthquake (Figs S4b and S5). Therefore, E2 was likely a new landslide (e.g. Yamada *et al.* 2012). We investigate the E2 event with a similar CSF modelling procedure





**Figure 3.** Waveform comparison between intermediate- and short-period data. (a) Self-normalized bandpass filtered (20–50 s) F-net waveforms aligned by the epicentre of the Tenryu landslide (E1). The yellow line shows the reference wavefront travelling at the phase velocity of  $3.11 \text{ km s}^{-1}$ . (b) Self-normalized bandpass filtered (0.125–0.5 s) Hi-net waveforms aligned by the epicentre of the Tenryu landslide (E1). The instrumental response of the Hi-net data was corrected. The yellow line shows the reference wavefront travelling at the phase velocity of  $3.11 \text{ km s}^{-1}$ . (c) Self-normalized bandpass filtered (0.125–0.5 s) Hi-net waveforms aligned by the epicentre of the  $M_w$  5.1 earthquake determined by the GCMT project at  $36.730^\circ\text{N}$ ,  $140.530^\circ\text{E}$  on 2011-03-18 23:49:34 (UTC). Blue and red lines are the predicted  $P$ - and  $S$ -wave arrivals (IASP91 model; Kennett & Engdahl 1991).

and find that the event can be well explained as centroid single forces (Fig. S6). The estimated duration is 24 s and the maximum centroid force is  $0.34 \times 10^{10} \text{ N}$ . Furthermore, we identify a third coherent phase  $\sim 3.5$  min before the signals associated with the E2 event (Figs S3 and S4a). The amplitudes of these signals are about 50 per cent of those of the E2 event and the signals are about 30-s long. Our surface wave detector cannot locate this seismic event (E3) due to the poor signal-to-noise ratios of records at further away stations. However, this event is likely close to the E2 event because the near-field stations at different azimuths recorded almost equal-separation times between the phases of E2 and E3 (Fig. S3). Although the signals from E3 are less clear in short-period records (Fig. S5), leaving their physical origins ambiguous, they are unlikely regular earthquakes due to a lack of body wave phases.

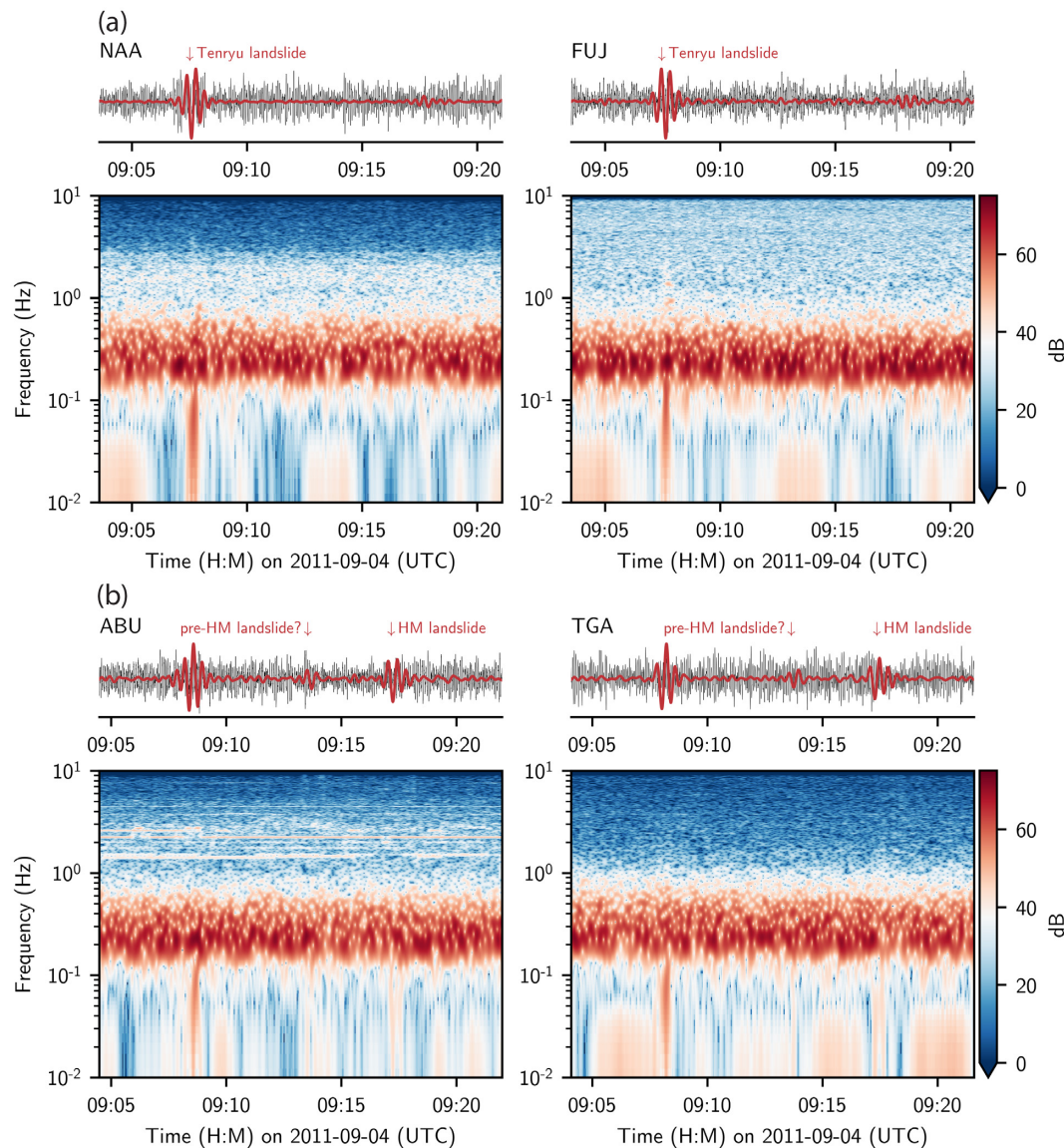
## 5 DISCUSSIONS

### 5.1 Source characteristics of E1, E2 and E3

The seismic events detected in this study (E1, E2 and E3) are unlikely typical earthquakes. The seismic sources generated signals that are distinctly different from those of regular earthquakes. For

regular earthquakes, for example, a moment magnitude ( $M_w$ ) 5.1 earthquake near E1 (with the source duration  $\sim 1$  s), seismic waveforms have clear  $P$ - and  $S$ -wave arrivals, and both short-period ground motions can be identified up to 300 km away (Fig. 3c). However, we do not observe clear  $P$  and  $S$  waves from the E1 at a similar distance range (Fig. 3b). The lack of clear short-period body waves makes it difficult to locate these sources with standard techniques. In contrast, we observe clear and coherent intermediate-period (20–50 s) surface waves at stations up to 3000 km (Fig. 2a). Although focusing and defocusing effects due to lateral structural heterogeneity can regulate surface-wave amplitudes in complex ways (e.g. Dalton & Ekström 2006), the lack of short-period body waves and the clear intermediate-period surface waves (Figs 3 and 4) collectively suggest these abnormal seismic radiations are not from typical earthquakes.

Seismic event E1 in Tenryu Ward, Shizuoka city is likely a landslide that was identified by the local forest office in Shizuoka prefecture. This landslide was reported 3 days after our resolved event time and is within 5 km of our detected seismic source location (Fig. 5e). The landslide was further confirmed by the aerial photos from the Geospatial Information Authority of Japan (Geospatial Information Authority of Japan 2011) and can be clearly identified in the optical

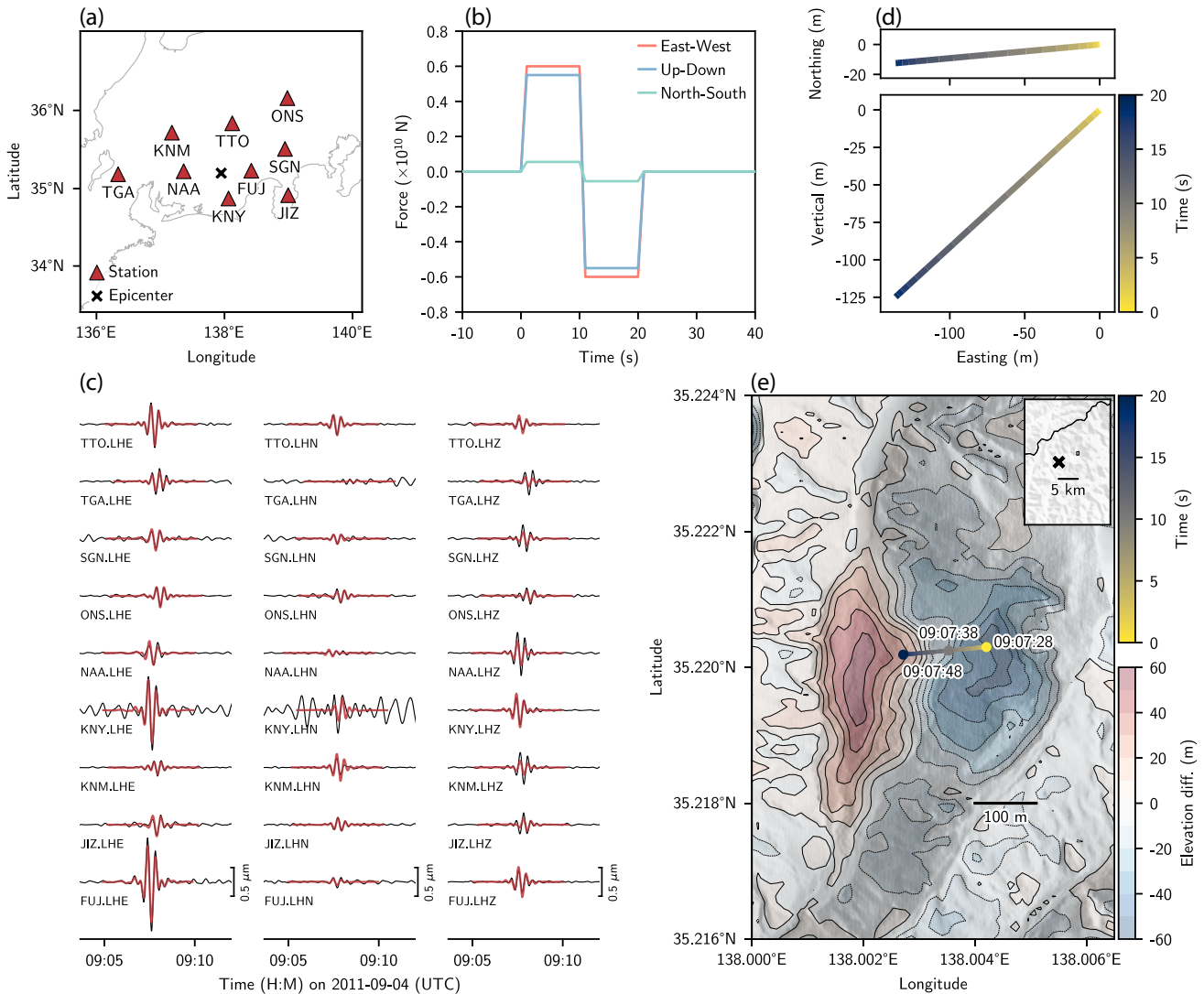


**Figure 4.** Spectrograms of the near-field broadband F-net records near (a) the Tenryu (E1) and (b) the Higashi-Matadani (HM, E2) landslides. Upper section of each panel shows the self-normalized filtered waveforms applying the 4th order Butterworth high-pass (100 s; black) and bandpass (20–50 s; red) filters. Bottom section is the spectrogram. The station code is denoted on the left-top of each panel. The station location is shown in Figs 5 a and S6a.

satellite imageries (Fig. 6). The field survey used a Laser Profiler to construct a digital elevation model (DEM). By differencing the DEMs before and after the landslide, the elevation changes show that the mass slid 200 to 500 m along the slope from east to west with a width of  $\sim 300$  m (Fig. 5e). The DEM model suggests that the Tenryu landslide displaced a total volume of  $1.2 - 1.5 \times 10^6$  m<sup>3</sup>, covering a region of  $\sim 9.0 \times 10^4$  m<sup>2</sup> with a maximum thickness of  $\sim 50$  m (Kanto Regional Forest Office Japan 2012; Seo *et al.* 2012; Yumoto & Takashima 2013). Assuming an average density of  $2.6 \times 10^3$  kg m<sup>-3</sup>, the landslide might have displaced a total mass of  $3.1$  to  $3.9 \times 10^9$  kg. To further evaluate the landslide dynamics, we compare the Rayleigh- to Love-wave amplitude ratios of the observed waveforms and the predicted ratio from our preferred CSF model (Fig. 7). The observations show an asymmetric radiation pattern of Rayleigh and Love waves, representing the sliding directivity from east to west, and the model-predicted pattern agrees well with the observations.

The two seismic events (E2 and E3) near Higashi-Matadani in Mie prefecture may relate to ones of the deep-seated landslides reported from field surveys after the typhoon transit. The E2 event is likely the Higashi-Matadani landslide, the largest field-reported landslide adjacent to the seismically determined location (Sakai 2011; Numamoto *et al.* 2012). Following a scaling relationship in Ekström & Stark (2013), we estimate the mass of the Higashi-Matadani event (E2) as  $1.8 \times 10^9$  kg and the volume as  $7.0 \times 10^5$  m<sup>3</sup> from the resolved CSF model, assuming a density of  $2.6 \times 10^3$  kg m<sup>-3</sup> (Yamada *et al.* 2013). The CSF model shows the mass displaced from south-east to north-west as a first order approximation of the mass movement. The model agrees well with the general changes of the topography, but cannot resolve the details, including possible redeposition that is modulated by the steep valley (Fig. 8). The third event (E3) occurred  $\sim 3.5$  min before the Higashi-Matadani landslide (E2) (Figs 4b, S3 and S4), but is challenging to locate with the current dataset. This event is likely the





**Figure 5.** Summary of the centroid single force (CSF) modelling and the digital elevation models (DEMs) of the Tenryu landslide (E1). (a) Distribution of the stations used for the CSF modelling. (b) The inverted three-component force–time function. (c) Black and red lines are the observed and synthetic waveforms, which are bandpass filtered at 20–50 s. Station codes and channels are listed on each column. (d) East–north and East–vertical trajectories (displacements) of the centre of mass. Colour represents the time. (e) Coloured contour denotes the differentiation of DEMs before and after the landslide. Coloured line is the trajectory of the centre of mass, along with the time on 4 September 2011 (UTC). The inset is the regional map. The small rectangle is the area of Fig. 5(e). The black line denotes the administrative boundary.

**Table 2.** Parameters of the CSF models for the Tenryu landslide (E1).

Duration (s)	Minimum misfit	$F_{\max}$ ( $\times 10^{10}$ N)	$F_{\text{Up-Down}}$ ( $\times 10^{10}$ N)	$F_{\text{North-South}}$ ( $\times 10^{10}$ N)	$F_{\text{East-West}}$ ( $\times 10^{10}$ N)
10	0.300	2.02	1.35	0.145	1.50
20	0.282	0.82	0.55	0.055	0.60
24	0.288	0.67	0.45	0.040	0.50
30	0.353	0.67	0.45	0.035	0.50
34	0.439	0.64	0.40	0.040	0.50
40	0.610	0.64	0.40	-0.060	0.50
50	0.738	0.65	0.40	0.080	0.50

Mochiyama-Tanigawa landslide, which is located about 1 km north-west of Higashi-Matadani landslide (Sakai 2011; Numamoto *et al.* 2012). If so, the source area of this landslide is about 30 per cent of the Higashi-Matadani landslide. However, the occurrence time of the Higashi-Matadani landslide is 40 min before our detection time

as reported by local residents (Numamoto *et al.* 2012). The timing inconsistency undermines the landslide hypothesis. However, no coherent seismic phases can be identified from near-field records 40 min before our detection. Alternatively, the smaller signals may come from a precursory event of the Higashi-Matadani landslide. We examine the 3-D particle motions of four near-source stations. The analysis shows similarities between the particle motions of the events E2 and E3. However, this qualitative exercise cannot conclusively determine the relation between E2 and E3, largely because of the noisy horizontal F-net records.

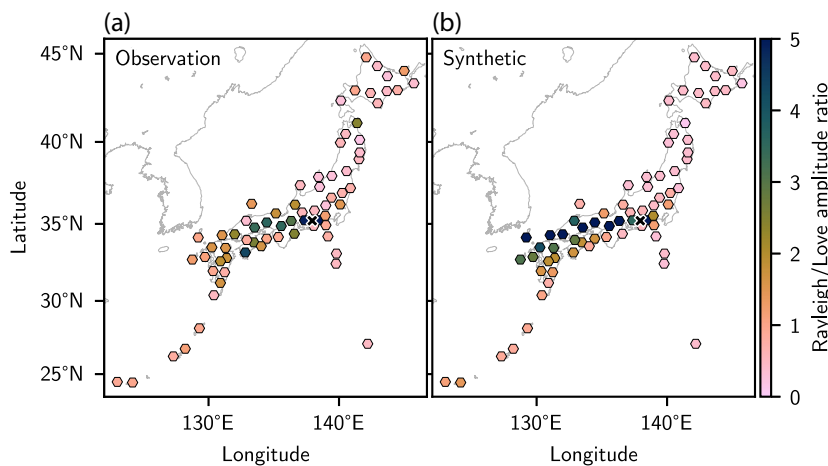
## 5.2 Empirical scaling relationship of attributes of Tenryu landslide E1

Our preferred CSF model of the Tenryu landslide (E1) has a maximum centroid force ( $F_{\max}$ ) of  $0.82 \times 10^{10}$  N, suggesting a total





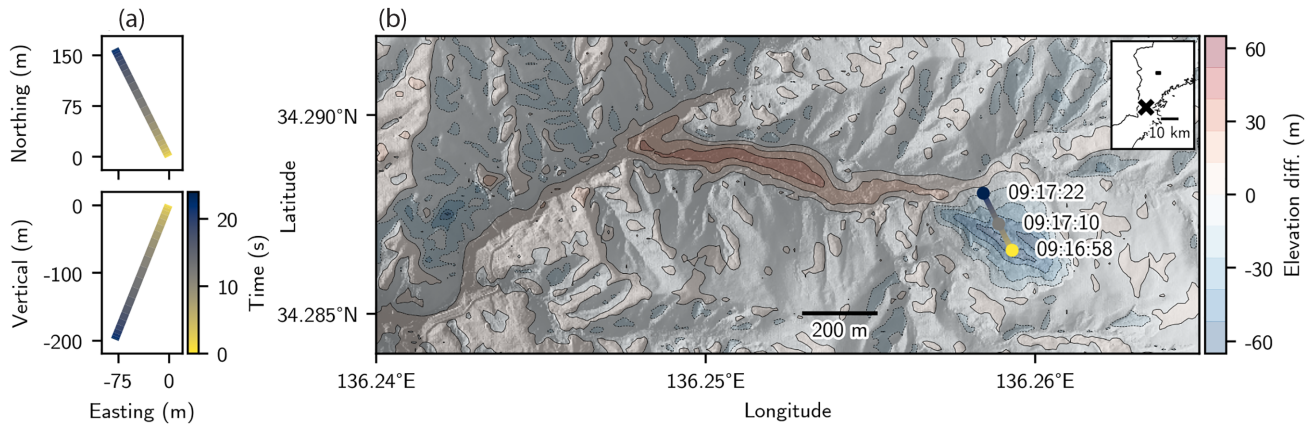
**Figure 6.** Google Earth™ imagery (a) before (13 May 2010) and (b) after (15 November 2011) the Tenryu landslide (E1), provided by Maxar Technologies.



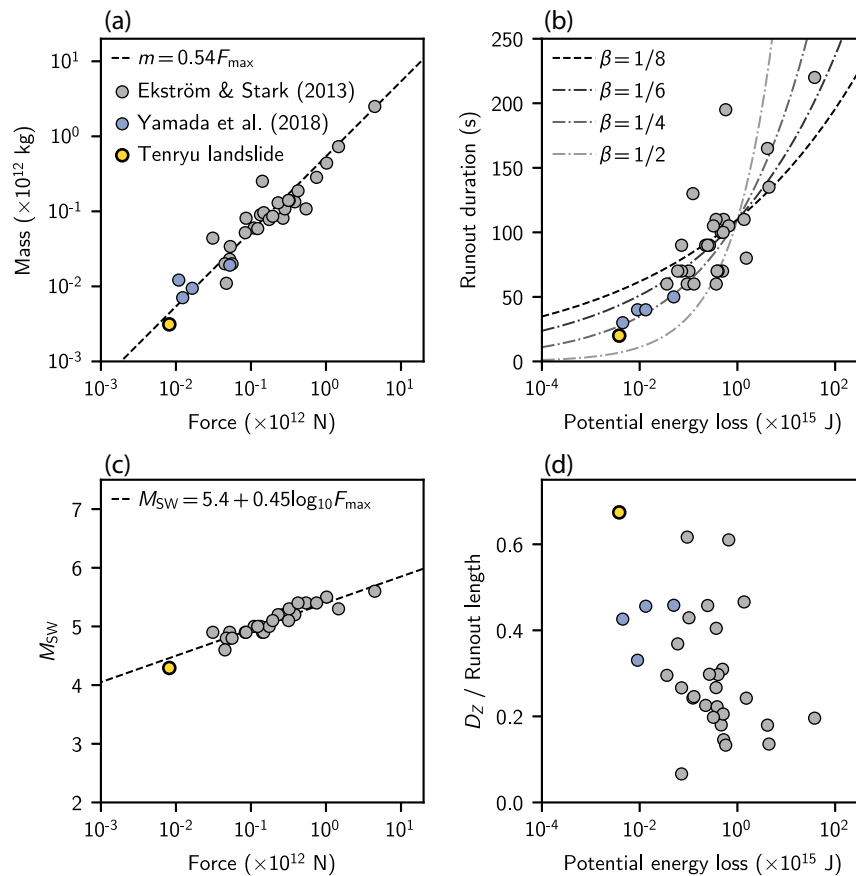
**Figure 7.** Comparison of the (a) observed and (b) synthetic Rayleigh- to Love-wave amplitude ratio for the F-net data. Each marker corresponds to a F-net station, colour coded by its Rayleigh- to Love-wave amplitude ratio. The cross marker shows the epicentre of the Tenryu landslide (E1). The synthetic Rayleigh- and Love-wave are computed from our preferred CSF model of the Tenryu landslide (E1).

displaced mass of  $4.4 \times 10^9$  kg when assuming an empirical scaling relationship in Ekström & Stark (2013). To understand the landslide dynamics, we explore the CSF model uncertainties by examining an ensemble of models that can explain the observations within 5 per cent of the minimum misfit ( $\leq 0.296$ , Table 2, Fig. S7). This exercise suggests that the  $F_{\max}$  is likely within  $0.77 \pm 0.06 \times 10^{10}$  N, indicating that the displaced mass ranges from  $3.8$  to  $4.5 \times 10^9$  kg.

This seismically inferred total mass agrees with the field survey estimate, despite that the empirical scaling relationship was drawn from landslides ten times larger than the Tenryu event (Fig. 9a). For example, the Siachen landslides in the high mountains of Pakistani Kashmir deposited mass complexes on the order of  $0.188 \times 10^{12}$  kg and generated centroid forces on the order of  $\times 10^{11}$  N (Ekström & Stark 2013). However, the total masses ( $m$ ) and the centroid peak



**Figure 8.** Mass trajectory for the Higashi-Matadani landslide (E2) and the topography change of the Higashi-Matadani landslide site. (a) East–north and east–vertical trajectories (displacements) of the centre of mass. Colour represents the time. (b) Coloured contour denotes the differentiation of digital elevation models (DEMs) before and after the landslide. Background topography is the DEM after the landslide. Coloured line is the trajectory of the centre of mass, along with the time on 4 September 2011 (UTC). The inset is the regional map. The cross marker is the epicentre of the Higashi-Matadani landslide. The small rectangle is the area of Fig. 8(b). The black line denotes the coastline and the administrative boundary.



**Figure 9.** Scatter plot of landslide parameters. (a) Maximum centroid force ( $F_{\max}$ ) versus landslide mass. The Tenryu landslide mass in this study is from field observations. (b) Potential energy loss  $\Delta E$  versus runout duration  $\Delta t$ . The curves plot  $\Delta t = 110 \Delta E^\beta$  (Ekström & Stark 2013, for  $\beta = 1/8$ ). (c)  $F_{\max}$  versus surface wave magnitude ( $M_{\text{SW}}$ ). (d) Potential energy loss versus the ratio of the vertical mass-centre displacement ( $D_z$ ) and runout length. The runout length corresponds to the summation of the east–west, north–south, up–down displacement vectors from the CSF modelling.

forces of both the Siachen landslides and the Tenryu landslide follow the same scaling relationship,  $m = 0.54 F_{\max}$  (Ekström & Stark 2013, Fig. 9a). Further, the maximum momentum from the CSF model and the  $M_{\text{SW}}$  magnitude of the Tenryu landslide fit other scaling relationships proposed in Ekström & Stark (2013) as well

(Fig. 9c). These agreements validate the scaling relationships over a large range of landslide sizes (Ekström & Stark 2013).

With the seismically estimated mass, we can further obtain the sliding acceleration history and the failure trajectory of the Tenryu landslide from the CSF model (Fig. 5d, Allstadt 2013; Gualtieri &

Ekström 2018), which agrees well with the field survey observations (Fig. 5e). The results show promises of using seismic observations to obtain accurate landslide trajectories in remote regions where satellite images or field surveys may be limited. To understand the landslide movement, we also estimate the dynamic frictional coefficient  $\mu$  with a total mass of  $3.11 \times 10^9$  kg (Text S1; Brodsky *et al.* 2003; Yamada *et al.* 2013), which ranges from 0.23 to 0.46 (Fig. S8), concurring with  $\mu$  of documented major landslides ( $0.2 \leq \mu \leq 0.6$ , e.g. Mt St Helens; Brodsky *et al.* 2003). The obtained frictional coefficient(s) is also proportional to the displaced volume ( $V$ ) as  $\mu \sim V^{-0.0774}$ , which scaling relationship suggests a possible velocity-weakening friction law that uniformly applies to small and large landslides (Lucas *et al.* 2014).

The potential energy loss  $\Delta E$  of the Tenryu landslide do not scale with its runout duration  $\Delta t$  (Fig. 9b) as  $\Delta t \propto \Delta E^{\frac{1}{8}}$  proposed in Ekström & Stark (2013). This is likely because the vertical displacement is comparable to the runout length of the Tenryu landslide (Fig. 5d), in contrast to the landslides dominated by horizontal movements in other regions. The Tenryu landslide occurred within a narrow valley and displaced along a steep slope, which is underlain by the alternated layers of sandstone and mudstone (Fig. 5e; Kanto Regional Forest Office Japan 2012; Yumoto & Takashima 2013). The layers are the Late Cretaceous accretionary-sedimentary rocks that develop fragile textures involving fractures and joints (Kanto Regional Forest Office Japan 2012). Similar geological predispositions of deep-seated landslides are also found in the southwest direction on the ridgeline of the landslide (Fig. 5e). High erosion rate due to the extreme climate and active tectonic regime may have facilitated the development of high-relief mountains and steep hills across the Japanese island, which likely causes landslides in the region with short durations and large vertical displacements (Oguchi *et al.* 2001; Yamada *et al.* 2018). The observed differences between the Tenryu landslide and other catastrophic landslides support the hypothesis that the power-law coefficient (e.g.  $\frac{1}{8}$ ) reflects the topographical variations, which has also been observed in the other field, laboratory and analytic studies (Hibert *et al.* 2011; Levy *et al.* 2015; Farin *et al.* 2018).

### 5.3 Outlook on real-time monitoring of landslides

The Tenryu landslide is  $\sim 400$  km east from the track of Typhoon Talas in a region with intense precipitations from the typhoon (Fig. 1b). Investigating such hazards away from the track requires a robust detection method that can effectively monitor a broad region. Our results suggest a useful detection algorithm that can identify small ( $\sim 100$  m scale) landslides with the existing F-net stations, and it is the first time the method being applied to detect and resolve previously unknown subaerial landslides in the region. Previous studies rely on a dense temporal network (Fan *et al.* 2020), and our results show promises to implement the technique to study environmental processes in regions with the existing permanent networks. Our approach is effective because it does not require phase-picking, prior knowledge of source types, or an accurate velocity model to calculate the traveltimes. Our approach uses local coherence across a triad, and stations of each triad would record waves travelling through a similar path. This helps to remove strong path effects of seismic wave propagation and hence is effective to detect remote landslides. Though requiring proper modifications of the method, our results suggest a future use of the Hi-net tiltmeter (e.g. Tonegawa *et al.* 2006) to systematically detect and locate landslides from 2004

to date. Such efforts will reveal the occurrence evolution and may aid deciphering the failure physics of the landslides in the region.

Although ground, aerial, and satellite methods can be used to map landslides with high spatial resolution, it is worth mentioning that it took 3 days for the local agencies to identify and survey the Tenryu landslide (Yumoto & Takashima 2013). These methods are often hampered by poor weather, restricting access and satellite visibility (e.g. Razak *et al.* 2013). In this study, we demonstrate that applying a suite of seismological analyses to regional seismic networks can effectively identify landslides from earthquakes and determine the dynamic processes of such failure events, including cross-examining sources resolved from our surface wave detector with standard catalogues, inspecting seismic wave signatures across multiple period bands, and modelling the failure histories as centroid single forces. Future combination of Hi-net ( $> 1$  Hz) and F-net (20–50 s) observations may provide an efficient way to screen our detected seismic sources as the frequency dependent radiation may be diagnostic to identify landslides. In conjunction with automatic classification algorithms (e.g. Dammeier *et al.* 2016; Manconi *et al.* 2016), our results show the possibility of using seismic records to resolve landslide locations and times in near-real time. Such data products can serve as preliminary results to assist future risk management and to guide rapid response of post-event surveys.

## 6 CONCLUSIONS

We detect and locate multiple landslides by applying the AELUMA method to 20-to-50-s period surface waves that were recorded by multiple spatially disconnected seismic arrays near Japan. These landslides occurred during the passage of Typhoon Talas 2011, including the Tenryu landslide (E1), the Higashi-Matadani landslide (E2) and the Mochiyama-Tanigawa landslide (E3). The Tenryu landslide displaced  $1.2 - 1.5 \times 10^6$  m<sup>3</sup> sediment and rock, and generated coherent intermediate-period Rayleigh waves that propagated up to 3000 km epicentral distance. Such signals are distinctly different from those of regular earthquakes. Our observations also show that landslide attributes, including the mass, inertial force and surface magnitude, empirically scale with each other and these scaling relationships are likely invariant for landslides of different sizes. Therefore, our methods are useful to identify small landslides and infer their physical attributes for regions with only sparse seismic networks. Our approach requires minimum assumptions and has potential to be implemented in near-real time for monitoring landslide activity and assisting future risk assessment.

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## DATA AVAILABILITY

Waveform data at F-net (<https://doi.org/10.17598/nied.0005>) and Hi-net (<https://doi.org/10.17598/nied.0003>) are available through NIED website (<https://www.hinet.bosai.go.jp/?LANG=en>). The facilities of IRIS Data Services and specifically the IRIS DMC (<https://ds.iris.edu/ds/nodes/dmc/>) are used for access to waveforms and related metadata. AELUMA MATLAB code bundle is available from IRIS DMC (<https://ds.iris.edu/ds/products/infrasound-aeluma/>). Green's functions used for the CSF modelling are provided by IRIS DMC Data Services Products: Synthetics Engine (<https://doi.org/10.17611/DP/SYNGINE.1>). The typhoon tracks are downloaded at [https://www.jma.go.jp/jma/jma-eng/jma-center/rsmc-hp-pub-eg/RSMC\\_HP.htm](https://www.jma.go.jp/jma/jma-eng/jma-center/rsmc-hp-pub-eg/RSMC_HP.htm). The AMeDAS precipitation data are downloaded at <https://www.data.jma.go.jp/gmd/risk/obsdl/index.php> (only in Japanese). The DEM data are available at <https://fgd.gsi.go.jp/download/menu.php> (only in Japanese). ObsPy (Beyreuther *et al.* 2010, version 1.1.0; <https://doi.org/10.5281/zenodo.165135>), matplotlib (Hunter 2007, version 3.0.3; <https://doi.org/10.5281/zenodo.2577644>), and the Generic Mapping Tools (Wessel & Luis 2017, version 6.1; <http://doi.org/10.5281/zenodo.3924517>) were used to generate figures. The CVX package (Grant & Boyd 2008, 2014, <http://cvxr.com/cvx>) was used for solving the least-square problem in locating source. The DEM data after the Tenryu landslide was provided by Chubu Regional Development Bureau, Ministry of Land, Infrastructure, Transport and Tourism, Japan. The DEM data of the Higashi-Matadani landslide was provided by the Geospatial Information Authority of Japan.

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## SUPPORTING INFORMATION

Supplementary data are available at *GJI* online.

**Table S1.** A list of detection (3–4 September 2011).

**Figure S1.** Example triad-sub array detection of the Tenryu landslide (E1). (a) Bandpass filtered (20-to-50 s) waveforms, colour-coded with the stations shown in (b). The shaded area is a 600-s time window that is used for beamforming. (b) The arrows are the observed propagation direction (black) and the predicted propagation direction (great circle path) from the Tenryu landslide (E1). The triangles are the station locations along with the station codes. The thick and thin black lines show the triad networks. The black dot is a triad centroid.

**Figure S2.** Example of a discarded source, #9 (Table S1). (a) Record section of source #9 (Table S1). The traces are self-normalized and band-pass filtered at 20–50 s period-band. The yellow line shows the reference wavefront travelling at a phase velocity of  $3.0 \text{ km s}^{-1}$ . (b) The thick and thin triangles are the triad subarrays. The arrow is the observed arrival angle. The colour for each dot represents the observed arrival time. The thin line between the epicentre and the centroid of each triad is the great circle path. The source is discarded because the aligned signals by the epicentre do not converge at the origin (Fig. S2a). The records are noisy, and the surface wave detector located the source with detections from subarrays of a narrow azimuthal range. This leads to a mislocation of the event, which might be the Iya landslide (Yamada *et al.* 2012) (Table S1).

**Figure S3.** Waveform comparison of the near-field F-net data. (a) Observed waveforms near the Tenryu landslide (E1). Station distribution is shown in Fig. 5(a). (b) Observed waveforms near the Higashi-Matadani landslide (E2). Station distribution is shown in Fig. S6a. Left-hand panels (Original) show the filtered waveforms applying the 4th order Butterworth-highpass (100 s) filter. The periods for the 4th order Butterworth-bandpass filter are denoted at upper left on other panels. Station code is denoted on each column.

**Figure S4.** Detection summary of the Higashi-Matadani landslide (E2). (a) Self-normalized bandpass filtered (20–50 s) waveforms aligned by the epicentre of the detected source. The green line shows the reference wave front travelling at the phase velocity of  $3.50 \text{ km s}^{-1}$ . The brown line indicates the coherent wave packets travelled from the precursory event E3 before 210 s from the origin time of the Higashi-Matadani landslide (E2). Inset is an enlarged section of the signals associated with the event E3. (b) The thick and thin triangles are the triad subarrays used and unused for the detection. The arrow is the measured arrival angle. The colour for each dot (centroid of each triad) represents the measured arrival time. The thin line between the epicentre and the centroid of each triad is the great circle path. The blue ellipse denotes the estimated location uncertainty. Background topography and bathymetry are from the GEBCO 2019 Grid (GEBCO Bathymetric Compilation Group 2019). The star denotes the  $M_{\text{JMA}} 1.7$  epicentre at  $33:832^\circ\text{N}$ ,  $134:985^\circ\text{E}$  on 09:17:32 (UTC).

**Figure S5.** Near-field short-period data for the Higashi-Matadani landslide (E2). (a) Hi-net station distribution. (b) Bandpass filtered (0.125–0.5 s) waveforms observed at Hi-net stations aligned by the location of the Higashi-Matadani landslide ( $34:287^\circ\text{N}$ ,  $136:259^\circ\text{E}$ , Numamoto *et al.* 2012). Closely overlapping stations (HYSH and URSH) are discarded for the visibility of the figure. The green line shows the reference wavefront travelling at the phase velocity of  $3.50 \text{ km s}^{-1}$  from the Higashi-Matadani landslide. The brown line is the reference wavefront travelling at the phase velocity of  $3.50 \text{ km s}^{-1}$  from the Higashi-Matadani landslide, but shifted – 210 s, indicating the possible signals of the Mochiyama-Tanigawa or the pre-Higashi-Matadani landslide. Blue and red markers are the predicted *P*- and *S*-wave arrival times (IASP91 model, Kennett &

Engdahl 1991) of the  $M_{\text{JMA}} 1.7$  earthquake ( $33:832^\circ\text{N}$ ,  $134:985^\circ\text{E}$ , 09:17:32 UTC).

**Figure S6.** Summary of the CSF modelling for the Higashi-Matadani landslide (E2). (a) Distribution of the stations used for the CSF modelling. (b) The three-components of the force-time function. (c) Black and red lines are the bandpass (20–50 s) filtered observed and synthetic waveforms. Station code and channel are denoted on each column.

**Figure S7.** Uncertainty of the CSF modelling for the Tenryu landslide (E1). Gray lines are the inverted force–time functions from the models that can produce similar misfits to the preferred solution, which are within 5 per cent of the minimum misfit ( $\leq 0.296$ ) for (a) east–west, (b) up–down and (c) north–south components. Coloured line is an optimum solution that gives the minimum misfit (Table 2). (d) Histogram of the maximum centroid force ( $F_{\text{max}}$ ) every  $0.05 \times 10^{10} \text{ N}$  bin from the models within 5 per cent of the minimum misfit ( $\leq 0.296$ ). Shaded area shows one standard deviation (SD). The mean and the one standard deviation of  $F_{\text{max}}$  for the suite of models are  $0.77 \pm 0.06 \times 10^{10} \text{ N}$ .

**Figure S8.** Estimates of dynamic frictional coefficients for the Tenryu landslide (E1). Digital elevation model (DEM) (a) before and (b) after the landslide. (c) Coloured contour denotes the differentiation of digital elevation models (DEMs) before and after the landslide. Background topography is the DEM after the landslide. (d) Topography profiles used for estimating the slope angles. Each profile corresponds to a line in Fig. S8(a). (e) Relationship between the slope angle and the dynamic frictional coefficient. The dynamic frictional coefficient  $\mu(\theta) = \tan \theta - F_h / (M g \cos^2 \theta)$  (Yamada *et al.* 2013), where  $\theta$  is the slope angle,  $F_h$  is the horizontal force ( $0.6 \times 10^{10} \text{ N}$ ) from the CSF model (Fig. 5b, Table 2),  $M$  is the mass ( $3.1 \times 10^9 \text{ kg}$ ) from the field survey and  $g$  is the gravitational acceleration. Shaded area is the range of estimates of  $\mu$  for the other catastrophic landslides Brodsky *et al.* (2003).

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