Rapid Seismic Waveform Modeling and Inversion With Neural Operators

Yan Yang[®], Angela F. Gao, Kamyar Azizzadenesheli[®], Robert W. Clayton[®], and Zachary E. Ross[®]

Abstract—Seismic waveform modeling is a powerful tool for determining earth structure models and unraveling earthquake rupture processes, but it is usually computationally expensive. We introduce a scheme to vastly accelerate these calculations with a recently developed machine learning paradigm called the neural operator. Once trained, these models can simulate a full wavefield at negligible cost. We use a U-shaped neural operator to learn a general solution operator to the 2-D elastic wave equation from an ensemble of numerical simulations performed with random velocity models and source locations. We show that full-waveform modeling with neural operators is nearly two orders of magnitude faster than conventional numerical methods, and more importantly, the trained model enables accurate simulation for velocity models, source locations, and mesh discretization distinctly different from the training dataset. The method also enables convenient full-waveform inversion with automatic differentiation.

Index Terms—Full-waveform inversion, geophysics, machine learning, partial differential equations (PDEs), waveform modeling.

I. INTRODUCTION

THE seismic wave equation relates displacement fields to external forces and the density and elastic structure in the Earth. Solutions to the wave equation form the basis of ground-shaking simulations of large earthquakes [1], [2], [3] and full-waveform inversion for Earth's structure [4], [5], [6]. Due to the highly heterogeneous nature of the Earth, as exemplified by subduction zones and sedimentary basins, there are no exact analytical solutions for these wavefields. Instead, approximate solutions are made possible by approximating derivatives through discretized spatial and time or frequency domains. Finite-difference methods (FDMs) have been popular since the early 1980 s due to their relatively straightforward formulation [7], [8], [9]. The spectral-element method (SEM), a particular case of finite-element methods (FEMs), which was

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introduced to seismology in the early 2000 s, combined the flexibility of FEMs with the accuracy of spectral approaches [10], [11], [12], [13]. These numerical solvers impose a trade-off between resolution and computation speed, with the computational cost proportional to the fourth power of frequency [14]. Thus, the cost of wave simulation is a major barrier to using full-waveform techniques for seismic inversion and updating models of the subsurface with new data.

A number of machine learning-based methods have been proposed in the past few years to provide a faster alternative for tackling seismological problems, such as signal denoising [15], [16], [17], event detection [18], [19], [20], and phase association [21], [22]. Deep neural networks have also recently been used to solve partial differential equations (PDEs), such as the Eikonal equation and wave equation [23], [24], [25], [26], [27]. These approaches to solve PDEs offer not only speedup in computational capabilities, but also low-memory overhead, differentiability, and on-demand solutions. Such advantages facilitate deep learning being used for seismic inversion [28], [29], [30], [31], [32]. However, one major limitation of these approaches is that the solutions generated by these models are dependent on the specific spatial and temporal discretization in the numerical simulation training set.

Recently, a paradigm named "neural operator" was developed to address the mesh-dependent shortcoming of classical neural networks by creating a single deep learning model that can be applied to different discretizations [33], [34], [35], [36]. This is made possible because neural operators can provably learn mappings between infinite-dimensional function spaces [37] and, therefore, are suitable for learning general solution operators to PDEs, which are valid even when the PDE coefficients (e.g., elastic properties) are varied. Since first introduced [33], a variety of neural operator models have been developed. In particular, the Fourier neural operator (FNO) is a model that uses the fast Fourier transform as an integral operator, and has been shown to outperform other neural operators in terms of efficiency and accuracy [38]. The FNO has been applied to many types of scientific problems, including weather forecasting [39], CO₂ sequestration [40], and coastal flooding [41].

Within the domain of seismology, neural operators were also recently used to learn general solution operators to the 2-D acoustic wave equation, a simplified case of the elastic wave equation [42]. This pilot study demonstrated that it was possible for a single FNO model to predict a complete wave-

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field given an arbitrary velocity model and mesh discretization. The success of this limited case highlights the potential of these methods; however, extending the method from the acoustic wave equation to the elastodynamic case requires substantially increased model complexity. By comparison with neural networks, FNO is not considered to be a deep architecture and is most analogous to the fully connected neural networks employed heavily until the 2010s. A U-shaped neural operator (U-NO) was recently proposed to enable very deep neural operators and facilitate fast training, data efficiency, and hyperparameter selection robustness [43].

In this article, we apply the U-NO architecture to full-seismic-waveform modeling. We train a U-NO model to learn a general solution operator to the 2-D elastic wave equation and demonstrate that the trained model enables fast and accurate simulation for source locations, velocity structures, and mesh discretization beyond the training dataset. The trained U-NO also allows for efficient full-waveform inversion with automatic differentiation.

II. METHODS

A. Neural Operator Learning

Operators are maps between function spaces, and the purpose of operator learning is to learn the operator given a dataset of input—output pairs. In seismology, it is common to write solutions to the wave equation, U(x), in terms of a linear integral operator acting on a source function, A(x)

$$U(x) = \int G(x, y)A(y)dy \tag{1}$$

where $x \in \mathbb{R}^4$ is the physical domain and G is a so-called Green's function defined for a particular velocity model. Equation (1) holds so long as the velocity model is not varied, because the wave equation remains a linear operator.

Instead, if we consider the case where the input function, A(x), is a velocity model, the solution operator, L, relating this to U(x) is nonlinear and cannot be written in the form of (1)

$$U(x) = (\mathcal{L}A)(x).$$

The most general version of the nonlinear solution operator L for the elastic wave equation is not known in closed form.

Neural operators are a class of models that aim to solve this problem, as they provably can learn a wide array of nonlinear operators. Their basic form consists of a composition of linear operators with nonlinear activations. More specifically, a neural operator with L layers can be written as follows:

$$v_0(x) = (PA)(x)$$

$$v_{l+1}(x) = \sigma(W_l v_l(x) + \int \kappa_l(x, y) v_l(y) dy), \quad l = 0, \dots, L-1$$

$$U(x) = (Qv_L)(x)$$
(2)

where v_l is the input function at the l^{th} layer, P is a pointwise operator that lifts the input function to a higher dimensionality, Q is a pointwise operator that projects the function back to the desired output dimensionality, W_l is a linear pointwise transformation that can keep track of nonperiodic boundary behavior, σ is a pointwise nonlinear activation operator, and

 κ_l is a kernel function that acts along with the integral as a global linear operator.

A neural operator is parameterized by P, Q, W_l , and κ_l . A critical aspect of this class of models is that these parameters are independent of the numerical discretization of the physical domain; i.e., they are shared across all possible discretizations in a similar way that in convolutional networks, the parameters are shared across neurons. It is this property that allows for the learning of maps between infinite-dimensional function spaces, as the discretization can be chosen dynamically at inference time independently of what was used for training.

If we are given a dataset of N numerical simulations, $\{A_i, U_i\}_{i=1}^N$, where the values of A_i are chosen to span the range of the expected function space, we can train a neural operator in a supervised fashion to map from arbitrary A into U.

Due to the expense of evaluating integral operators, neural operators may lack the efficiency of convolutional or recurrent neural networks in finite-dimensional settings. The FNO was proposed to mitigate this difficulty through the fast Fourier transform [38]. The kernel integral operator in (2) can be considered a convolution operator, defined in Fourier space as follows:

$$\int \kappa_l(x, y) v_l(y) dy = \mathcal{F}^{-1}(\mathcal{F}(\kappa_l) \cdot \mathcal{F}(v_l))$$
 (3)

where \mathcal{F} and \mathcal{F}^{-1} denote Fourier transform and its inverse, respectively. However, FNO imposes that each layer is a map between functions spaces with identical domain spaces, which may cause a large memory usage. The U-NO, an analogy to the U-net architectures, was proposed to allow progressively transforming the input function space with respect to a sequence of varying domains [43], [44]. After the lifting operator P, a sequence of L_1 nonlinear integral operators G_i is applied to v_0 and maps the input to a set of functions with decreasing dimensional domain. Then, a sequence of L_2 nonlinear integral operators G_i is applied to v_{L_1+1} and maps the input to a set of functions with increasing dimensional domain before the projection operator Q. Skip connections [44] are included to add vector-wise concatenation of v_{L_1+i} and v_{L_1-i} . The contracting and expanding parts are symmetric. The architecture of the U-NO used in this study is illustrated in Fig. 1, and we refer the interested readers to [33], [34], [35], [36], [38], and [43] for more details.

B. Numerical Simulation

We set up a training dataset of random source locations, S-wave velocity (V_S) models, and P- to S-wave velocity ratios (V_P/V_S). We define the velocity model on a 64 \times 64 mesh with 0.16-km grid spacing. The source is set as an isotropic explosive source randomly distributed on the mesh. V_S has an average background of 3 km/s and perturbed by random fields with a von Kármán covariance function with the following parameters: Hurst exponent $\kappa = 0.5$, correlation length $a_x = a_y = 8$ grids, and the fractional magnitude of the fluctuation $\varepsilon = 10\%$ background velocity. The power spectral density function of the von Kármán-type random field follows a power law (fractal randomness) and can accurately

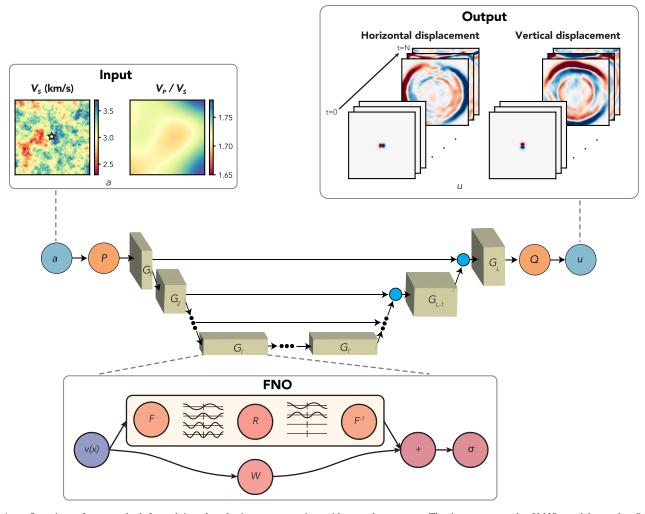


Fig. 1. Overview of our method for solving the elastic wave equation with neural operators. The inputs, a, to the U-NO model are the P- and S-wave velocity (V_P and V_S) models and the source location (indicated by the white star). (Top left) V_P is calculated from V_S and V_P/V_S , examples of which are shown. (Top right) Outputs, u, are the horizontal and vertical displacements at each time step, examples of which are shown. (Middle) U-NO architecture, orange circles P and Q denote point-wise operators, rectangles P denote general operators, and smaller blue circles denote concatenations in function space. (Bottom) Architecture of each FNO layer, where v is the input of the layer, P and P-1 are Fourier transform and its inverse, respectively, P and P-1 are nonlinear activation function. The expansion (or contraction) factors in (5) are set as follows: $c_{1,2}^s = \frac{3}{4}$, $c_{3,4}^s = \frac{1}{2}$, $c_{5,6}^s = 2$, $c_{7,8}^s = \frac{4}{3}$, $c_{1,2,3,4}^t = \frac{1}{2}$, $c_{5,6}^s = 2$, $c_{7,8}^t = \frac{4}{3}$, $c_{1,2,3,4}^t = \frac{1}{2}$, $c_{5,6}^t = 2$, and $c_{5,6,7,8}^t = \frac{1}{2}$.

represent the distribution of Earth's heterogeneity [45]. V_P/V_S is simplified to an average background of 1.732 perturbed by a smooth Gaussian random field with the following parameters: correlation length $\lambda = 32$ grids and standard deviation $\sigma =$ 2% background. This work, as our very first experiment to evaluate the feasibility of solving 2-D elastic wave equations, wants to focus on the parameters that the wavefield is most sensitive to. Therefore, we use the empirical relation between density and V_S to compute the density [46]. Other input parameters, such as density and attenuation, may be explored in future work. A total of 20 000 random sets of models are generated, and each of them is input to a GPU-based 2-D finite difference code in Cartesian coordinates to simulate the 2-D displacement field [47]. For simulation, the top boundary is set with a free-surface boundary condition, and the other three edges have absorbing boundary conditions. A total of 4s wavefield with a time step of 0.01 s and a major frequency content up to 6 Hz is simulated. Each simulation takes about 1.23 s with a GPU memory usage of 0.3 GB.

C. U-NO Model Training

We developed a framework that applies U-NO to the 2-D elastic wave equation. The architecture is depicted schematically in Fig. 1. U-NO takes the source location and V_P and V_S as inputs, where V_P is calculated from V_S and V_P/V_S . V_P and V_S are then passed through a point-wise lifting operator. A sequence of nonlinear integral operators (encoders) are applied that gradually contract the physical domain size after each inverse Fourier transform step, while simultaneously increasing the number of channels in the co-domain. These operators are followed by a sequence of nonlinear integral operators (decoders) that progressively expand the physical domain and decrease the number of channels. Finally, a pointwise projection operator leads to the output function [43]. The output of the U-NO model is the complete horizontal and vertical displacement wavefield functions over the medium domain, which can be queried at any mesh points desired, regardless of the input and output training mesh used.

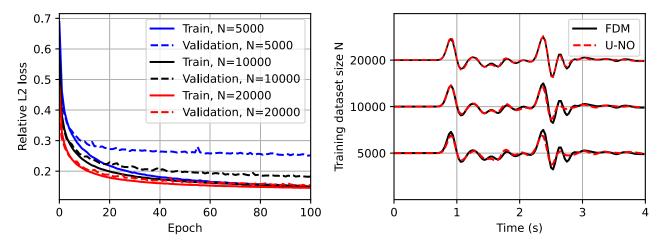


Fig. 2. Model performance as a function of the number of training samples. (Left) Relative L2 loss curves for the training and validation data. (Right) Example of simulated waveform comparison between FDM (black solid lines) and U-NO (red dashed lines). N means the number of simulations in the training dataset, including 90% for training and 10% for validation.

We describe the detailed parameters used in U-NO below following the notations in Fig. 1. The goal is to learn an operator mapping from the input function a to the output function u. The training is on an input mesh of $X_{\rm in} \times Y_{\rm in} \times T_{\rm in}$ and an output mesh of $X_{\rm out} \times Y_{\rm out} \times T_{\rm out} \times C_{\rm out}$, where $X_{\rm in} = Y_{\rm in} = X_{\rm out} = Y_{\rm out} = 64$, $T_{\rm in} = 3$ representing source, V_S , and V_P/V_S distribution on the mesh, $T_{\rm out} = 128$ for 32-Hz data output, and $C_{\rm out} = 2$ representing two displacement components (horizontal and vertical). This work applies the U-NO architecture designed for mapping between 3-D spatiotemporal function domains (x, y, t) without any recurrent composition in time [43]. The fourth dimension $C_{\rm out}$ of the output function u can be created in the last step through the projection operator Q. Constructing the operator to learn the mapping between 3-D spatiotemporal function domains

$$G: \left\{ a: [0, 1]^2 \times [0, T_{\text{in}}] \to \mathbb{R}^{d_A} \right\} \\ \to \left\{ u: [0, 1]^2 \times [0, T_{\text{out}}] \to \mathbb{R}^{d_U} \right\}. \tag{4}$$

The operators $\{G_i\}_{i=0}^L$, as shown in Fig. 1, that are used to construct the U-NO are defined as follows:

$$G_{i}:\left\{v_{i}:\left[0,\alpha_{i}\right]^{2}\times\mathbb{T}_{i}\rightarrow\mathbb{R}^{d_{vi}}\right\}$$

$$\rightarrow\left\{v_{i+1}:\left[0,c_{i}^{s}\alpha_{i}\right]^{2}\times c_{i}^{t}\mathbb{T}_{i}\rightarrow\mathbb{R}^{c_{i}^{c}d_{vi}}\right\}$$
(5)

where $[0, \alpha_i]^2 \times T_i$ is the domain of the input function v_i to the operator G_i , and c_i^s , c_i^t , and c_i^c are the expansion or contraction factors for the spatial domain, temporal domain, and co-domain for the ith operator, respectively. Note that $\mathcal{T}_0 = [0, T_{\rm in}], \ \alpha_0 = \alpha_{L+1} = 1, \ \text{and} \ \mathcal{T}_{L+1} = [0, T_{\rm out}].$ In this work, we set the number of layers to L=8. The details of the expansion and contraction factors c_i^s , c_i^t , and c_i^c are in Fig. 1. The lifting operator P to convert the input to a higher dimension channel space is a fully connected neural network with channel number $d_0=16$. The projection operator Q to the output domain is also a fully connected neural network. The activation function used in each FNO block is the Gaussian error linear unit (GELU) [48].

With the simulation dataset and the U-NO design, we train the U-NO model in a supervised manner with the objective of learning the general solution operator to the wave equation for arbitrary inputs. We divide the training dataset into 90% training and 10% validation. The model is trained with a batch size of 8. After hyperparameter tuning, the loss function we use in model training is the 90% relative L1 loss plus 10% relative L2 loss. The incorporation of L1-norm loss is more resistant to outliers. We use an Adam optimizer [49] with a learning rate of 10^{-3} and a weight decay of 10^{-5} . We trained for 100 epochs, which takes approximately 40 min/epoch using a single NVIDIA Tesla V100 GPU with 24-GB memory usage. A 70% of loss decrease is achieved in the first ten epochs. Once the U-NO model is trained, the model parameters require GPU usage of 3.8 GB, and the time for an evaluation on a new source and velocity model takes only 0.02 s with GPU usage of 0.9 GB.

III. RESULTS

A. Number of Simulations Needed for Training

Once completely trained, the U-NO model can be evaluated on a new input with very little computational cost (0.02 s compared with the FDM runtime of 1.23 s). The number of training simulations is the main factor in the computational cost. In the training process, we split the entire training dataset to 90% for training and 10% for validation. We test the performance of the model on the velocity models out of the training dataset. We can see that the U-NO model trained on a dataset of 5000 simulations can already predict the major phase arrivals, while increasing the dataset size from 5000 simulations to 20 000 provides better fit to the amplitudes (Fig. 2). With a training dataset of 20 000 simulations, the validation and training loss are very close, indicating there is no overfitting of the training data.

B. Generalizability to Arbitrary Velocity Structure or Discretization

The U-NO model is trained on random velocity models generated with the von Karman correlation function, which can best mimic the Earth's heterogeneous velocity distribution [45], [50]. We show by example that the U-NO model,

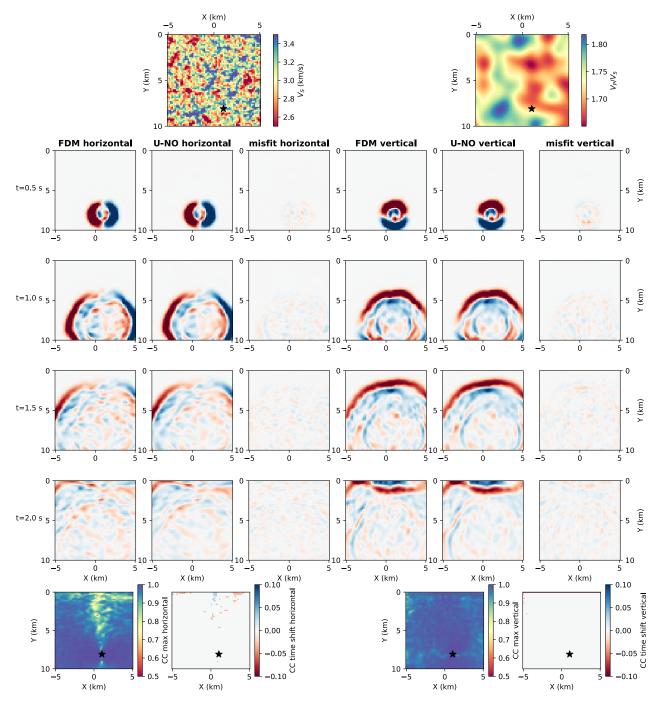


Fig. 3. Model generalization experiments 1: random fields of V_S and V_P/V_S model with four-times roughness of the training data. The top row shows V_S and V_P/V_S . From the second row to the fifth row, the wavefield snapshots at 0.5–2.0 s are shown. From left to right, the first three columns show the horizontal displacement of the FD simulation, U-NO prediction, and their misfit in the same color scale. The latter three columns show the vertical component. The horizontal and vertical displacement waveforms at each grid are cross correlated between FD and U-NO, with the maximum cross-correlation value and its associated time shift shown in the bottom row. For this case, the relative L2 loss of the U-NO simulation is 0.182.

although trained on random velocity models with some certain parameters, is applicable to arbitrary velocity models. These outcomes are, in fact, expected from theoretical grounds, because most physical functions can be approximated to arbitrary accuracy by random fields.

Our first example is with velocity models from a von Karman-type random distribution, but with a different covariance function than the one used for the training data. We increase the roughness of the velocity structure by a factor of 4 by decreasing the correlation length of V_S and V_P/V_S to only one-fourth that of the training data. As shown in Fig. 3, the wavefield snapshot has more coda than with the smoother models because of the scattering from increased heterogeneity. However, the coda waves are well modeled by U-NO when compared with the ground-truth simulation by FDM.

The velocity models used in the training data do not have coherent structures with discontinuities as in the real Earth, but wavefields for such models can still be simulated with

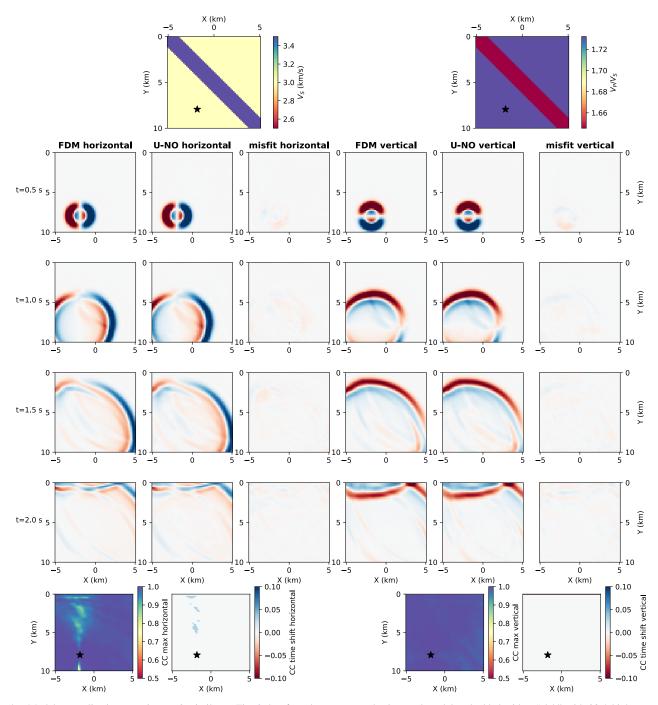


Fig. 4. Model generalization experiments 2: similar to Fig. 3, but for a homogeneous background model embedded with a "slab" with 20% higher V_S and 5% lower V_P/V_S . Relative L2 loss of the U-NO simulation is 0.090.

our method. As mentioned before, this is because discontinuous functions can be approximated to arbitrary accuracy by random fields. Fig. 4 shows a simple model with a dipping "slab" embedded in a homogeneous background. The slab has 20% higher V_S and 5% lower V_P/V_S . The wavefield snapshots show that the reflections from the high velocity anomaly are clearly predicted by U-NO. A more complex example is shown in Fig. 5, where a random subpanel of the Marmousi model, a 2-D velocity model with complex vertical and horizontal structures used in exploration studies [51], is used. The reflected and refracted waves are very complicated

due to the presence of folding and faulting, but the U-NO predictions still closely approximate the numerical solutions (Fig. 5).

One of the most important advantages of a neural operator compared with a neural network is its mesh-free nature, since it intrinsically learns the mapping between function spaces. A model trained on a particular mesh can be evaluated on any other mesh, even at finer spacing. The Fourier layers may learn from and evaluate functions on any discretization, because parameters are directly learned in Fourier space, and resolving the functions in physical space is simply projecting on the

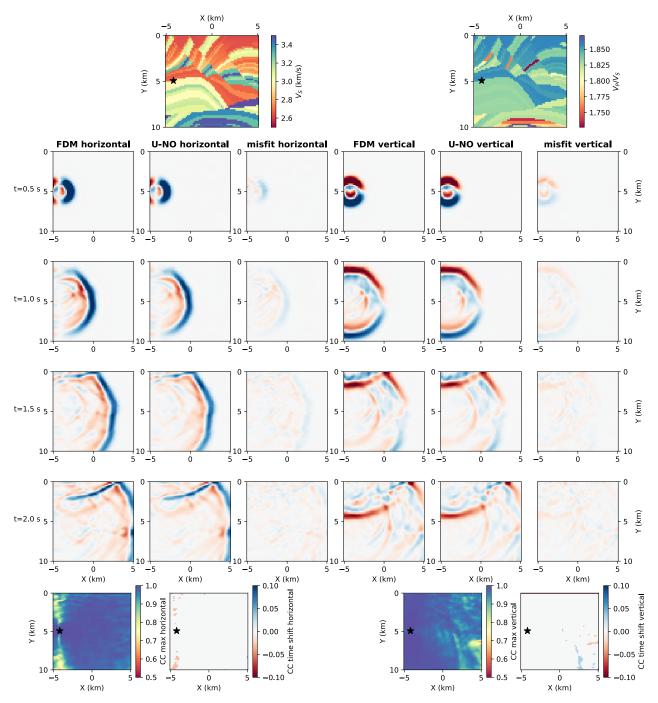


Fig. 5. Model generalization experiments 3: similar to Fig. 3, but for a random subpanel from the Marmousi model. The velocity perturbation range is normalized to 30% of the average velocity. Relative L2 loss of the U-NO simulation is 0.225.

basis [38]. The example in Fig. 6 shows the U-NO trained on a grid of 64×64 nodes applied to an input velocity model with 160×160 nodes. Here, both the input velocity model and the output wavefield can be seen at a much higher resolution, yet U-NO provides comparable prediction with the FDM solver. Note that if the resolution is increased by a factor of 2, a grid-based numerical solver, such as FDM, takes about six times greater computational time; however, the evaluation using U-NO takes only about 2.5 times longer, providing additional computational efficiency.

We evaluate the overall generalization performance of the trained U-NO by performing one thousand random realizations on each of these cases. The distribution of the relative L2-norm misfit and cross-correlation coefficient are plotted in Fig. 7. In the case of the Marmousi model, the extended tail of the histogram is attributed to the model's imbalanced complexity. In general, however, we see a very high cross-correlation coefficient (>0.95) between U-NO prediction and ground truth, confirming its robust generalizability.

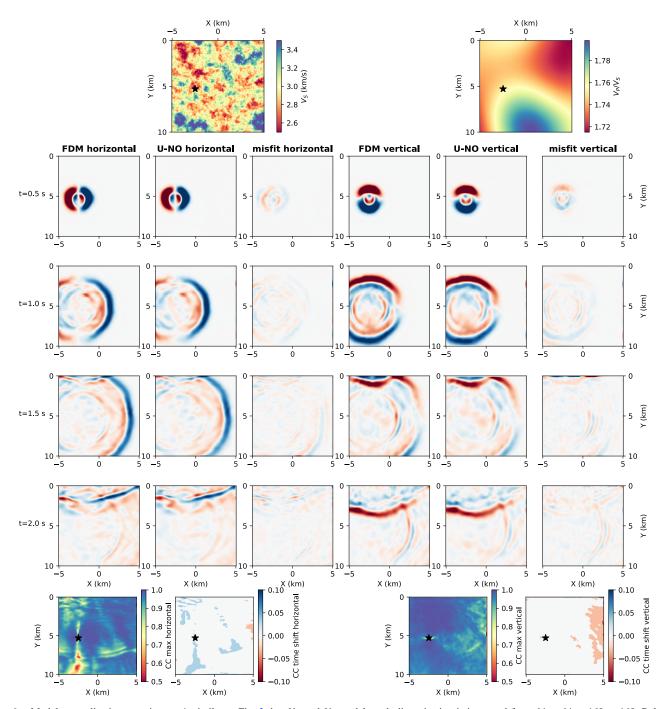


Fig. 6. Model generalization experiments 4: similar to Fig. 3, but V_P and V_S model mesh discretization is increased from 64×64 to 160×160 . Relative L2 loss of the U-NO simulation is 0.385.

C. Application to Full-Waveform Inversion

One of the most important applications of wavefield simulations is in full-waveform inversion (FWI), which uses the full recorded waveform to image the Earth's interior. The adjoint-state method is the traditional approach for computing the gradients of an objective function with respect to parameters of interest [4], [5]. Neural operators are differentiable by design, which enables gradient computation with reverse-mode automatic differentiation. It has been shown that automatic differentiation and the adjoint approach are mathematically equivalent [28]. Hence, the trained U-NO model allows for

convenient FWI, and the associated speed and accuracy should depend only on the forward modeling part.

We demonstrate the inversion performance using the velocity structure of random subpanels in the Marmousi model [51]. The synthetic waveform data are simulated with FDM [47] using 14 sources distributed in a ring shape. In Fig. 8, we use the true source location, receivers on all 64×64 grids, and noise-free waveform data; the goal here is not to demonstrate resolution, but rather the computational accuracy of the method. We then invert for V_P and V_S simultaneously by starting with homogeneous initial V_P and V_S models and forward propagating the wavefield with the

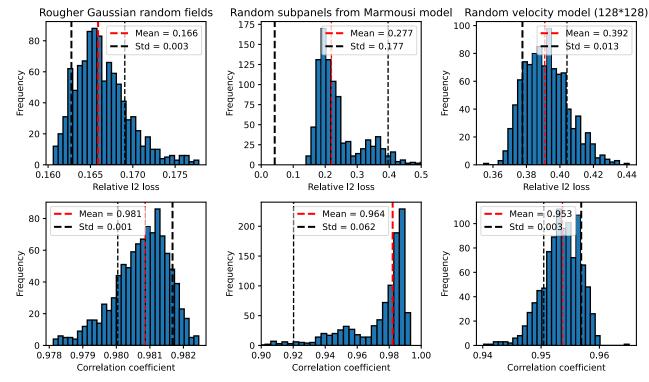


Fig. 7. Output evaluation. (Top) Distribution of relative L2 loss and (bottom) correlation coefficient between the U-NO predictions and ground truth. From left to right, the columns are corresponding to the experiments in Figs. 3, 5, and 6. The red and black dashed lines mark the mean and standard deviation of the histograms.

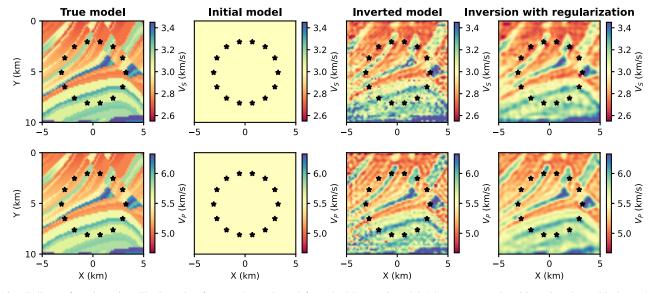


Fig. 8. Full-waveform inversion. The inversion for a random subpanel from the Marmousi model; 14 sources are placed in a ring shape (black stars), and receivers are placed at every node of the 64×64 grid. From left to right, the columns represent true velocity model, initial model for inversion, inverted model without regularization, and inverted model with the zeroth- and first-order Tikhonov regularization. The top and bottom rows are the models for V_S and V_P , respectively.

U-NO for each source. The misfit is defined by the meansquare error between the forward modeled and true wavefield. The gradient of the misfit with respect to V_P and V_S can be computed through automatic differentiation. V_P and V_S are then iteratively updated with gradient descent for 100 iterations using the Adam optimizer [49] with a learning rate of 0.01. Each iteration takes only about 1.4 s by taking advantage of U-NO forward computation. The results in Fig. 8 show a relative L2-norm misfit between the true and inverted model of only 3%. This successful inversion, in turn, further validates the accuracy of forward modeling with U-NO.

Besides the fact that the inversion target velocity model is quite different from the smooth random fields in the training dataset, this experiment itself is difficult due to conventional problems in full-waveform inversion, such as cycle skipping (multiple local maxima in the least-squares misfit function).

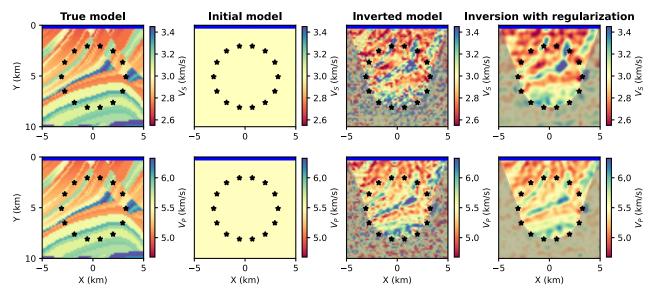


Fig. 9. Full-waveform inversion. Similar to Fig. 8, but the receivers are only placed on the 64 grids on the surface (blue line on the top of each subpanel). The gray shaded areas mask the areas without ray path coverage.

We also show that if we only use 64 receivers on the surface, the inversion results in the region with ray path coverage are still reasonably accurate (Fig. 9). Inversion with a biased homogenous initial model is also capable of producing relatively accurate results (see Supplementary Material).

IV. DISCUSSION AND CONCLUSION

We use the relative L2 loss between the FDM and U-NO predictions to evaluate the performance of the trained model for generalization. The relative L2 loss is defined as the L2-norm of the difference between the prediction and ground truth divided by the L2-norm of the ground truth. This ratio is used to evaluate the performance of the trained model. When using the same mesh discretization as the training data, the relative L2 loss is around 10%–20%, but this number rises to 30%–40% when the tests are performed on finer grids (Fig. 7). These values are misleading, however, because the relative L2 loss imposes equal weights to the entire sparse matrix of waveforms that is dominated by small amplitudes close to 0. Alternatively, the cross-correlation coefficient is a quantity that is more sensitive to the seismic phases with amplitudes larger than background noise. A cross-correlation coefficient larger than 0.95 suggests that the coherence of the U-NO prediction is excellent, even for the scenarios with large relative L2 loss (Fig. 7). In addition, the FWI results confirm that the large L2loss is not so important, since even challenging models can still be properly recovered (Fig. 8).

Besides the more than an order of magnitude higher speed, the most important advantage of the neural operator-based full-waveform modeling is its generalizability to arbitrary velocity models or discretization. This is because the neural operator learns a general solution operator to the wave equation instead of a specific instance of input velocity models. Once the neural operator is trained, it can be used by the entire seismology community for any region of a similar size without the need for retraining. Since the full-waveform modeling with a neural operator has easily accessible gradients for convenient FWI,

we anticipate that this approach will eventually make FWI as affordable as travel time tomography.

One of the main limitations of the method is the domain extent. For a trained neural operator, the function is defined on a fixed domain extent (e.g., it could be a unit cube). We can evaluate at a different grid size but cannot change the extent. We are now working on an extension of the work, where we recursively predict the wavefield. Through this way, a trained neural operator is essentially taking the first few time steps as input and then output the next few time steps, and there will be no need for retraining.

The scalability of evaluation using a trained neural operator with respect to the grid size and the number of time steps is a little different from conventional FDM. Assuming the original dimension is (Nx_1, Ny_1, Nt_1) , where Nx_1 and Ny_1 are the number of grids in the x and y domain, respectively, and Nt_1 is the number of time steps. If the new dimension is (Nx_2, Ny_2, Nt_2) , the memory becomes $(Nx_2 \cdot Ny_2 \cdot Nt_2)/(Nx_1 \cdot Nt_2)$ $Ny_1 \cdot Nt_1$) times the original memory, which is consistent with the FDM. In the example presented in this article, evaluation using U-NO takes three times the GPU memory of the FDM approach, and this scaling should be consistent with increasing grid points. In terms of computational cost, the majority of it for UNO is on the Fourier transform and its inverse. The computational cost of fast Fourier transform with dimension (Nx_1, Ny_1, Nt_1) is proportional to $Nx_1 \cdot Ny_1 \cdot Nt_1 \cdot$ $\log (Nx_1 \cdot Ny_1 \cdot Nt_1)$, and therefore, the new computational time becomes $(Nx_2 \cdot Ny_2 \cdot Nt_2 \cdot \log (Nx_2 \cdot Ny_2 \cdot Nt_2))/(Nx_1 \cdot Ny_2 \cdot Nt_2)$ $Ny_1 \cdot Nt_1 \cdot \log (Nx_1 \cdot Ny_1 \cdot Nt_1)$ times the original computational time. This scaling is slightly higher than that of FDM; however, considering the 60 times acceleration in the example presented in this article, UNO evaluation on an increased dimension of $1024 \times 1024 \times 1024$ should still have ≈ 40 times the acceleration.

The most compute- and memory-intensive part of the UNO method is the one-time training process. The cost of training for the 2-D case is tractable on a single GPU. For the extension

from 2-D to 3-D modeling, the computation and memory will increase due to the larger dataset and the larger number of parameters to learn. Therefore, the next step is to enhance data compression and parallelization to accelerate the training process and reduce the storage. Since this is a learning-based approach, the model performance can be improved by fine-tuning the model parameters and increasing the size of the training dataset. More importantly, any future advancements made in neural operator model architectures will be able to be directly incorporated into the system as they occur. For example, the improvement from linear layers of FNO to U-NO enables faster training convergence. As a result, we should only take current performance metrics as a starting point.

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